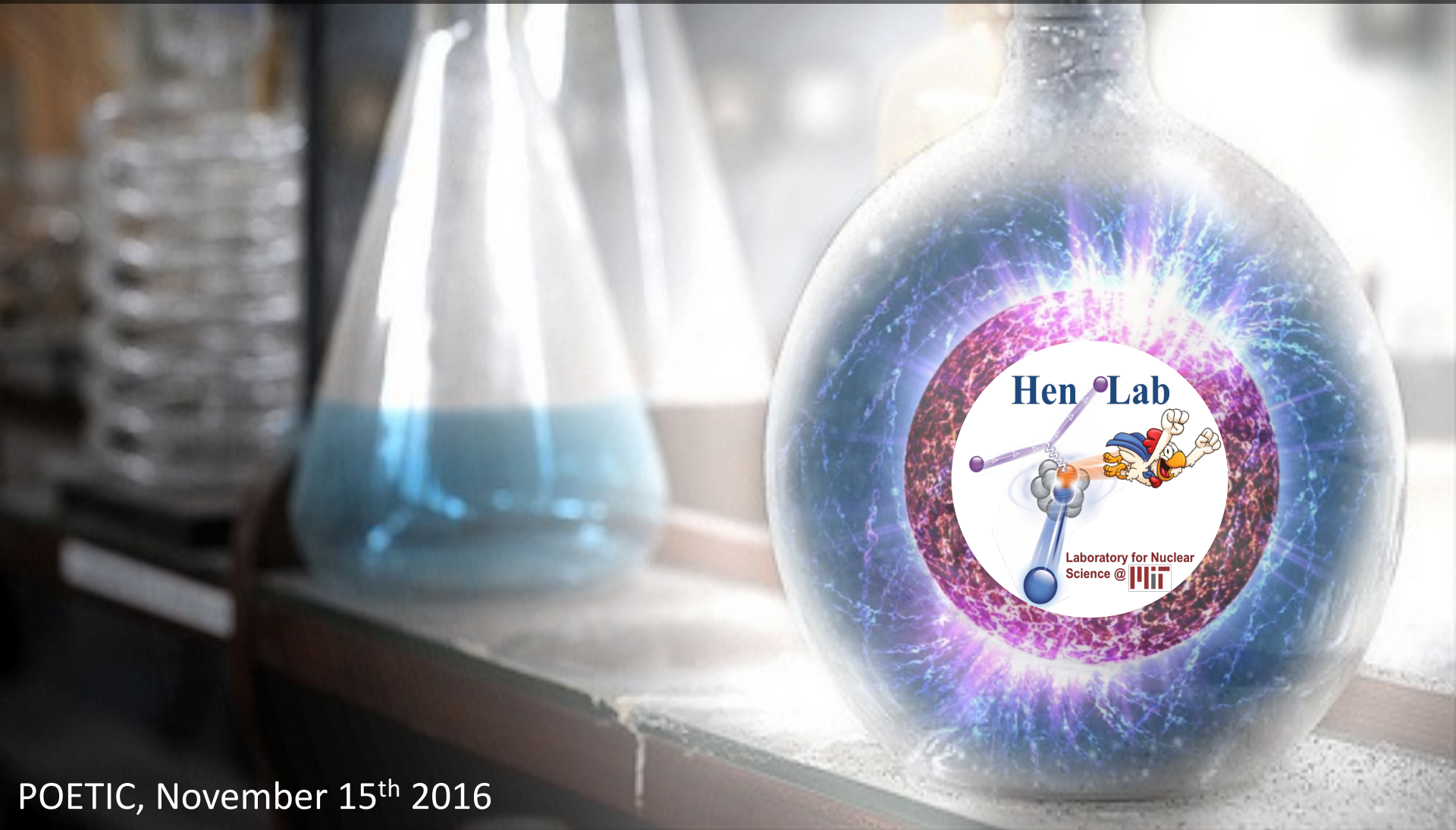


Short-Range Nuclear Structure

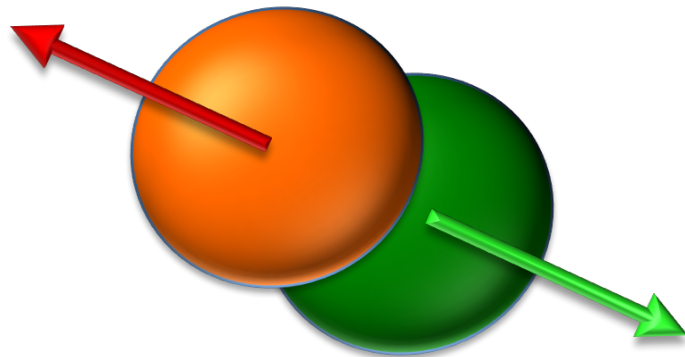
Or Hen – MIT





SRC are pairs of nucleon that are close together in the nucleus (wave functions overlap)

=> Momentum space: pairs with high relative momentum and low c.m. momentum compared to the Fermi momentum (k_F)



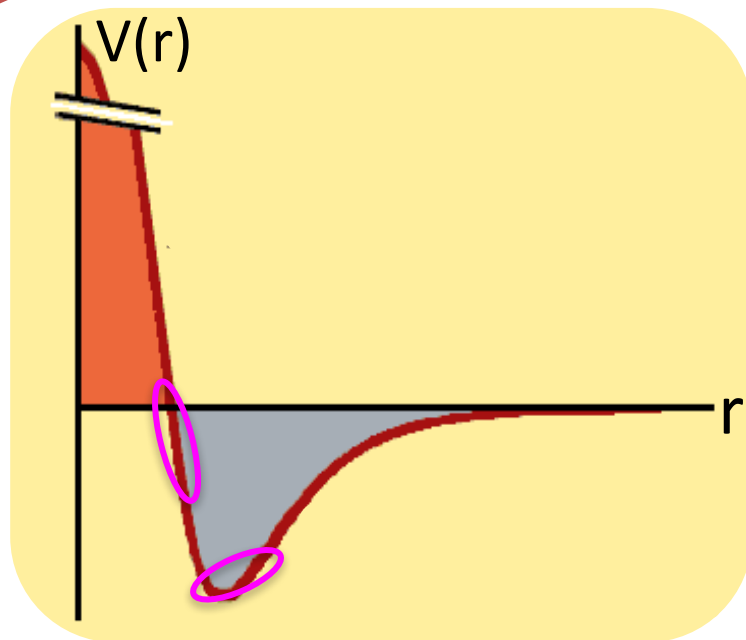


Why SRC?



Nuclear Physics

Better understanding of the
nucleon-nucleon interaction and the
nuclear momentum distribution



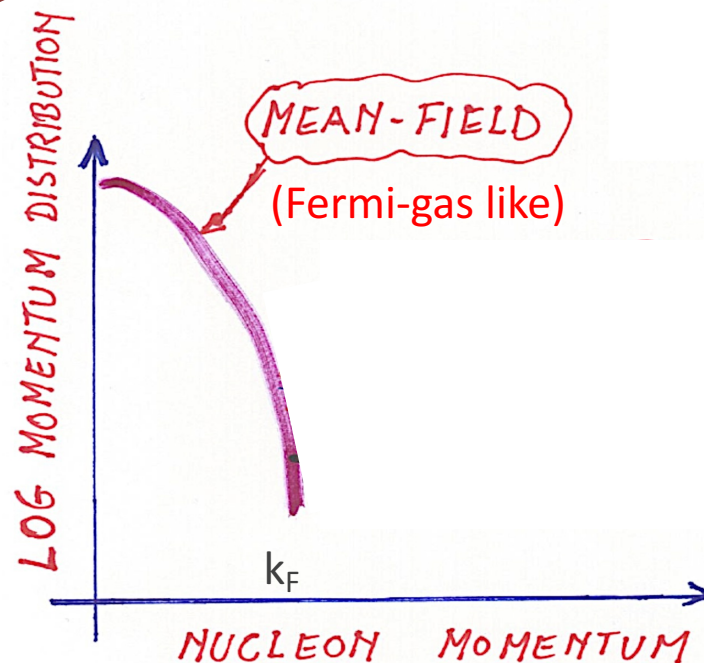
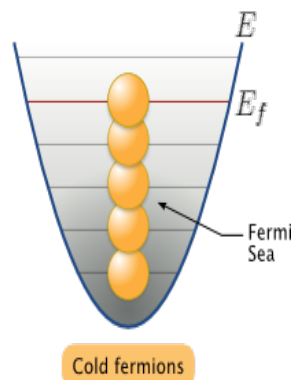


Why SRC?



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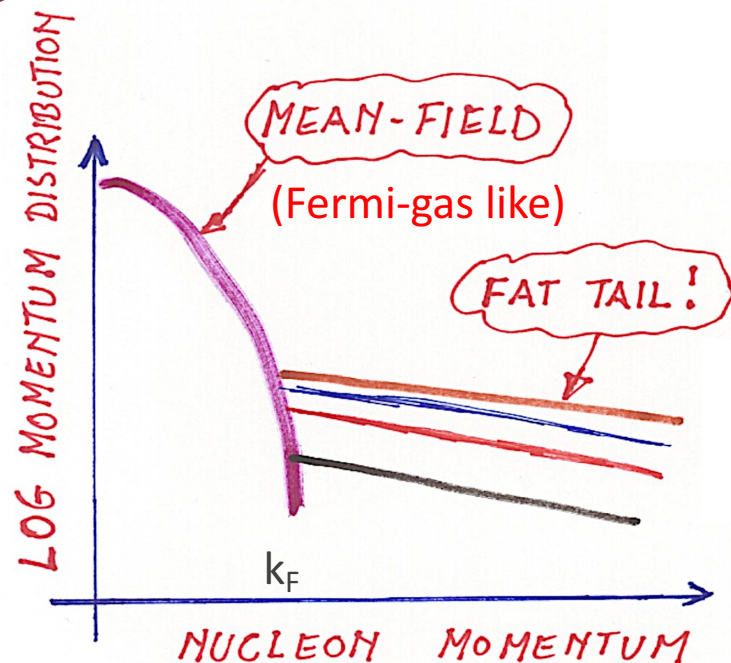
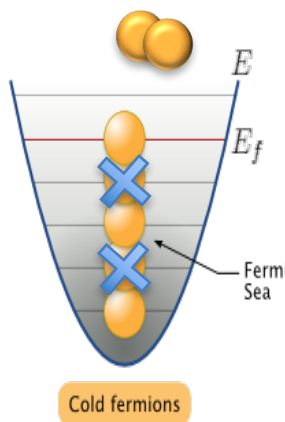


Why SRC?



Nuclear Physics

Better understanding of the
nucleon-nucleon interaction and the
nuclear momentum distribution





Why SRC?



You can't do nuclei without
correlations!



Why SRC?



Today: (short) overview of SRC and a presentation of a new effective theory for SRC in nuclei

You can't do nuclei without correlations!



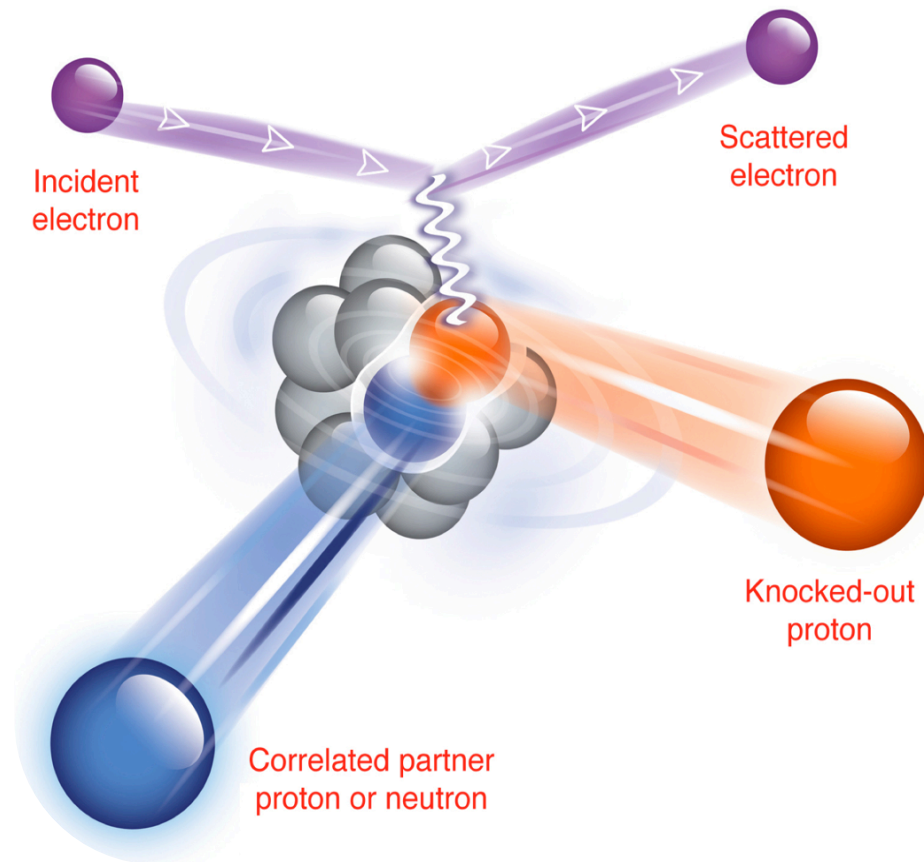
Exclusive 2N-SRC Studies



Breakup the pair =>

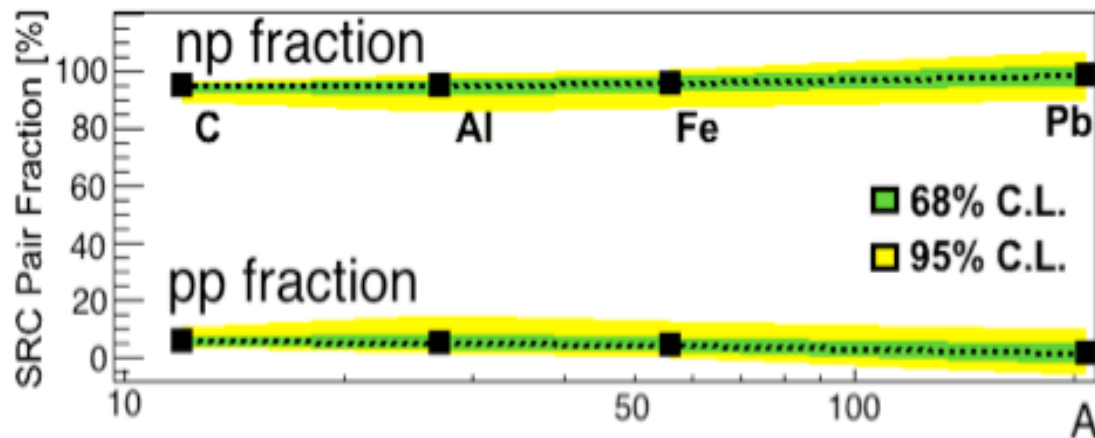
Detect both nucleons =>

Reconstruct 'initial' state

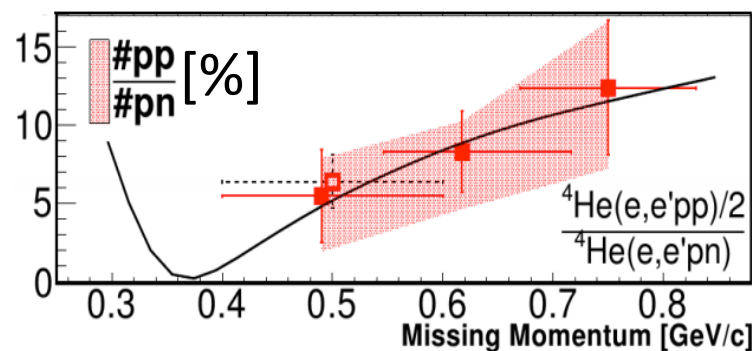




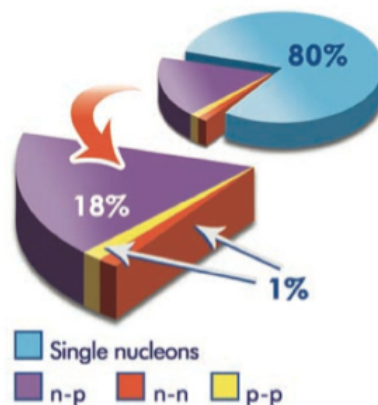
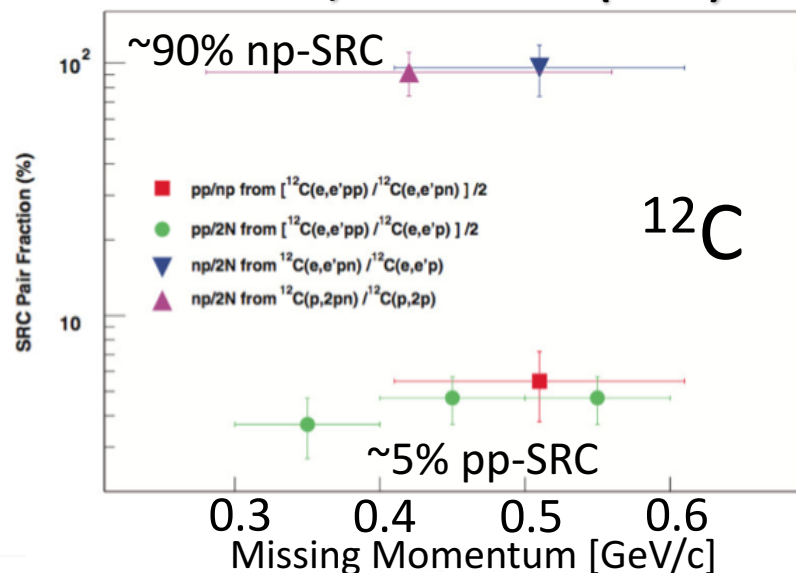
SRC Isospin Structure



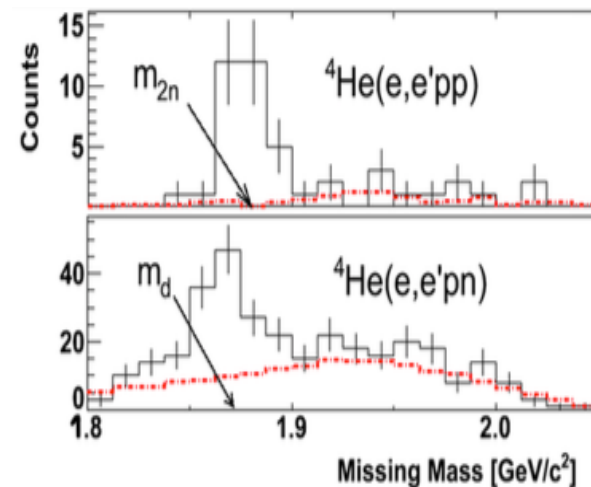
O. Hen et al., Science 364 (2014) 614



R. Subedi et al., Science 320 (2008) 1476



I. Korover et al., PRL 113 (2014) 022501



A. Tang et al., PRL (2003);

E. Piasezky et al., PRL (2006);

R. Shneor et al., PRL (2007)

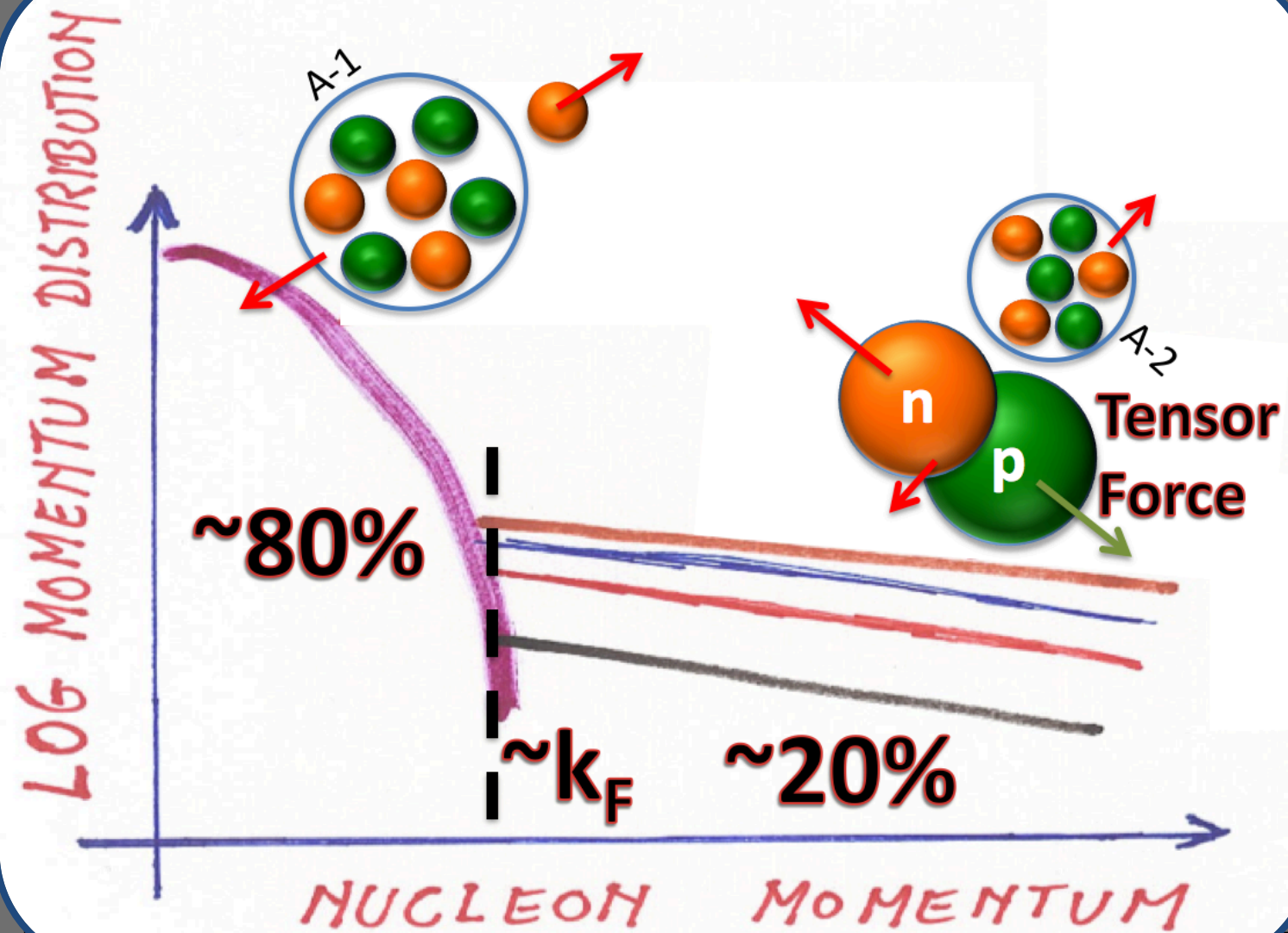


Bottom Line:

- SRCs account for:
 - ~ 20% of the nucleons in nuclei.
 - ~100% of the high-p ($k > k_F$) nucleons in nuclei.
- Predominantly due to np-SRC.
- Universal for $A = 4 - 208$ nuclei.
- Tensor force dominance at short distance.

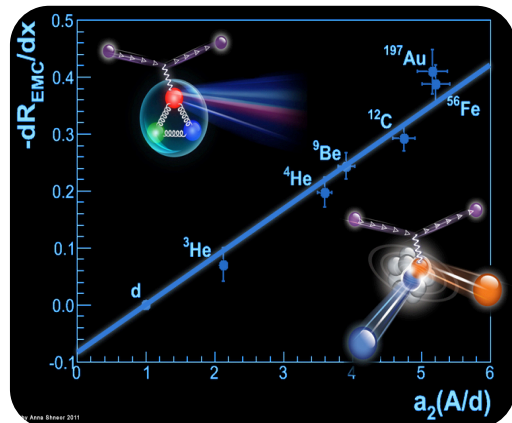
0.3 0.4 0.5 0.6
Missing Momentum [GeV/c]

Universal structure of nuclear momentum distributions

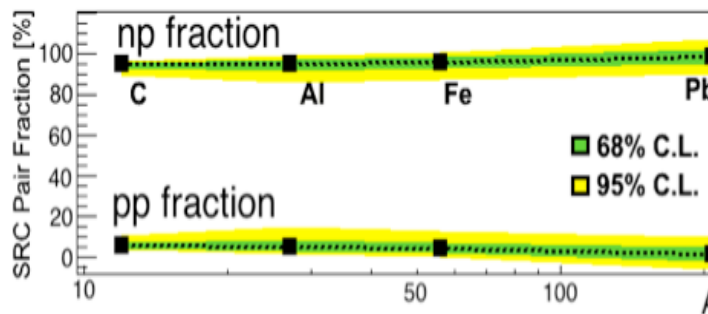




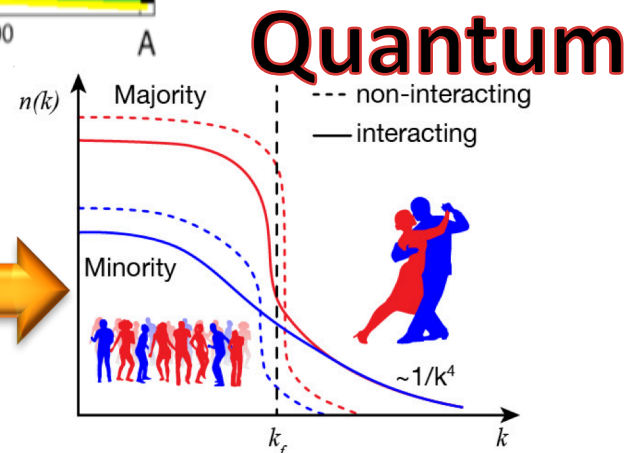
Importance of SRC Properties



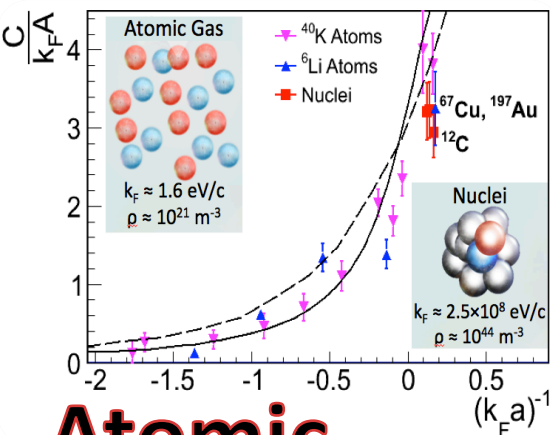
Particle



Nuclear

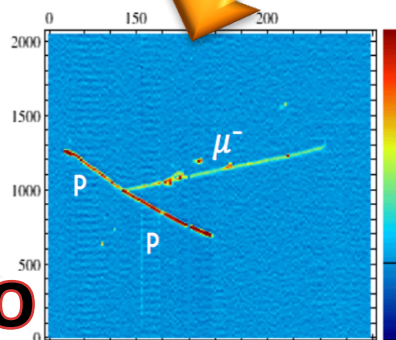


Quantum

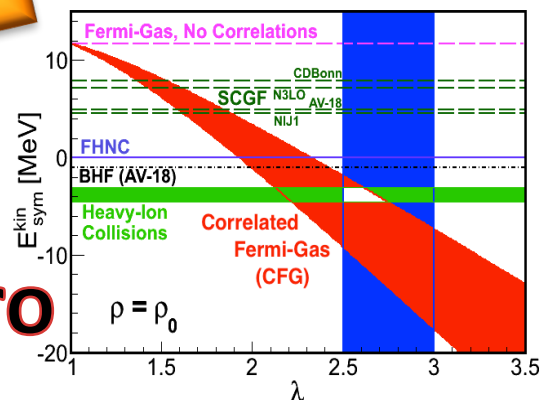


Atomic

Neutrino



Astro



Two-component interacting Fermi systems

The contact term

Please forget about nuclear
physics for a moment





The Contact and Universal Relations



Concept developed for dilute two-component Fermi systems with a short-range (δ -like) interaction.

$$\text{dilute} \equiv r_{\text{eff}} \ll a, d$$

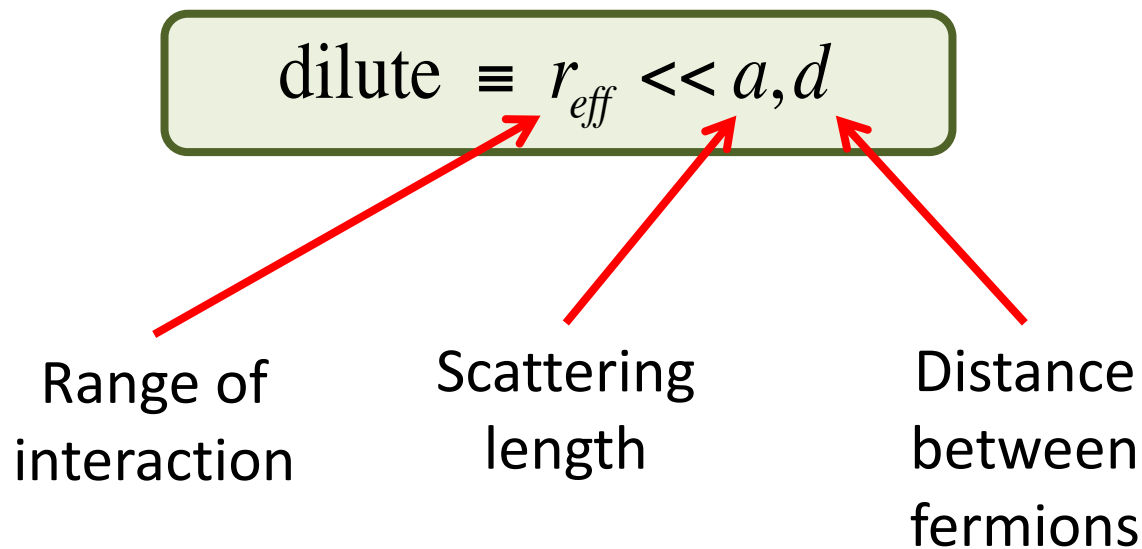
Dilute System



The Contact and Universal Relations



Concept developed for dilute two-component Fermi systems with a short-range (δ -like) interaction.



Range of interaction much smaller than the other relevant length scales in the problem

Dilute System



The Contact and Universal Relations



Contact interaction is represented through a boundary condition

Imposing this B.C. on the Schrödinger equation yields an asymptotic wave function when two fermions get very close

$$\Psi \xrightarrow{r_{ij} \rightarrow 0} \underbrace{(1/r_{ij} - 1/a)}_{\text{Two Body}} \underbrace{A_{ij}(\mathbf{R}_{ij}, \{\mathbf{r}_k\}_{k \neq i,j})}_{\text{A-2}}$$

Dilute System



Short Distance
Factorization



The Contact and Universal Relations

$$\Psi \xrightarrow{r_{ij} \rightarrow 0} (1/r_{ij} - 1/a) A_{ij}(\mathbf{R}_{ij}, \{\mathbf{r}_k\}_{k \neq i,j})$$



$$n(k) = C / k^4 \quad \text{for } k > k_F$$

Dilute System



Short Distance
Factorization



High
Momentum Tail



The Contact and Universal Relations



$$\Psi \xrightarrow{r_{ij} \rightarrow 0} \underbrace{(1/r_{ij} - 1/a)}_{\text{red bracket}} A_{ij}(\mathbf{R}_{ij}, \{\mathbf{r}_k\}_{k \neq i,j})$$



$$n(k) = C / k^4 \quad \text{for } k > k_F$$

Dilute System



Short Distance
Factorization



High
Momentum Tail



The Contact and Universal Relations

$$\Psi \xrightarrow{r_{ij} \rightarrow 0} (1/r_{ij} - 1/a) A_{ij}(\mathbf{R}_{ij}, \{\mathbf{r}_k\}_{k \neq i,j})$$



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Dilute System



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Factorization



High
Momentum Tail



The Contact and Universal Relations



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$$n(k) = C / k^4 \quad \text{for } k > k_F$$

Tan's Contact term:

1. Measures the number of SRC different fermion pairs.
2. Determines the thermodynamics through a series of universal relations.

Dilute System



Short Distance
Factorization



High
Momentum Tail

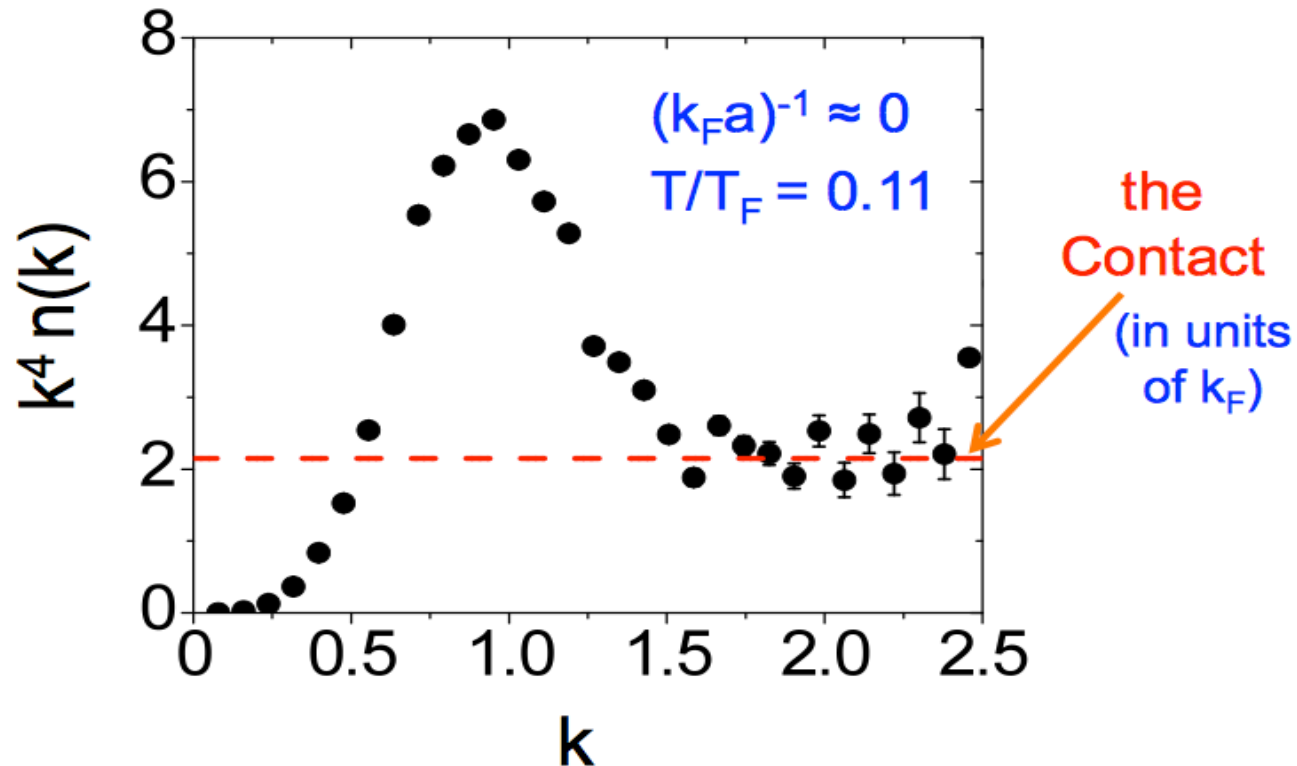


Experimental Validation



Two spin-state mixtures of ultra-cold ^{40}K and ^6Li atomic gas systems.

=> extracted the contact and verified the universal relations



Stewart et al. PRL **104**, 235301 (2010)



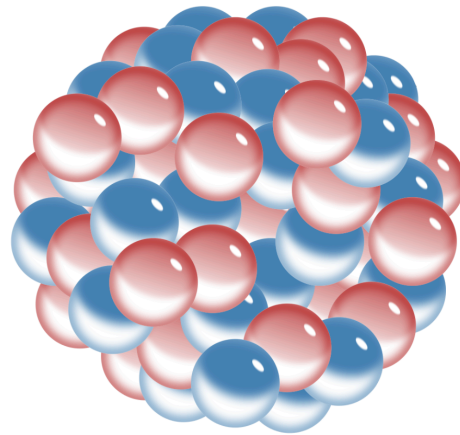
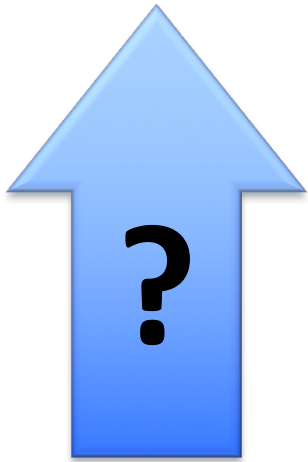
What About a *Nuclear* Contact ?



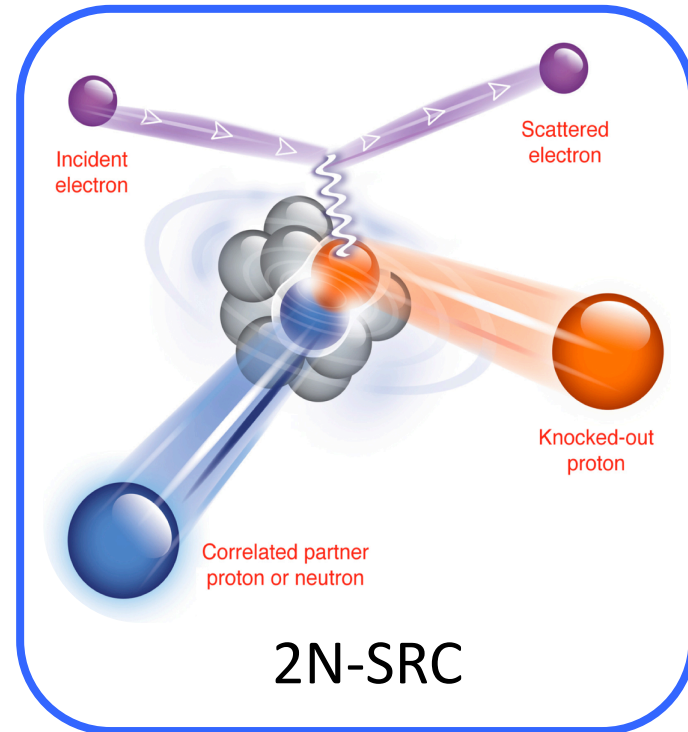
A Nuclear Contact?



Concept developed for:
dilute two-component Fermi systems with a short-range interaction.



protons and
neutrons





Theory Says: not so much



Are nuclei dilute? (i.e. $r_{\text{eff}} \ll a, d$)

$$r_{\text{eff}} \approx \frac{\hbar}{2 \cdot m_{\pi} \cdot c} \approx 0.7 \text{ fm}$$

[Tensor force]

$$d = \left(\frac{\rho}{2} \right)^{-1/3} \approx 2.3 \text{ fm}$$

$$a(^3S_1) = 5.42 \text{ fm}$$

$$r_{\text{eff}} (0.7 \text{ fm}) < d (2.3 \text{ fm}), a (5.4 \text{ fm})$$



But Experiment Says....



Is there $1/k^4$ scaling regardless?

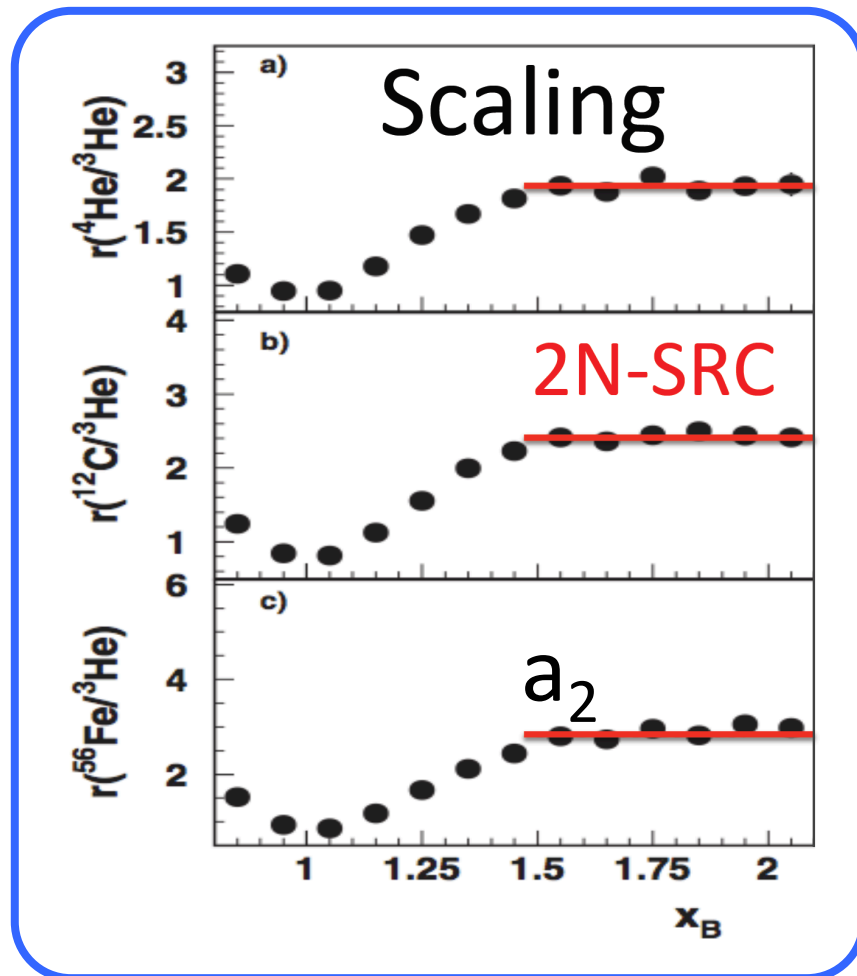
$$1.5k_F < k < 3k_F$$

$$n_A(k) = a_2 (A/d) \cdot n_d(k)$$

nucleus A
momentum
distribution

deuteron
momentum
distribution

experimental
constant



The momentum distribution of nucleons in medium to heavy nuclei is proportional to that of deuteron at high momenta.



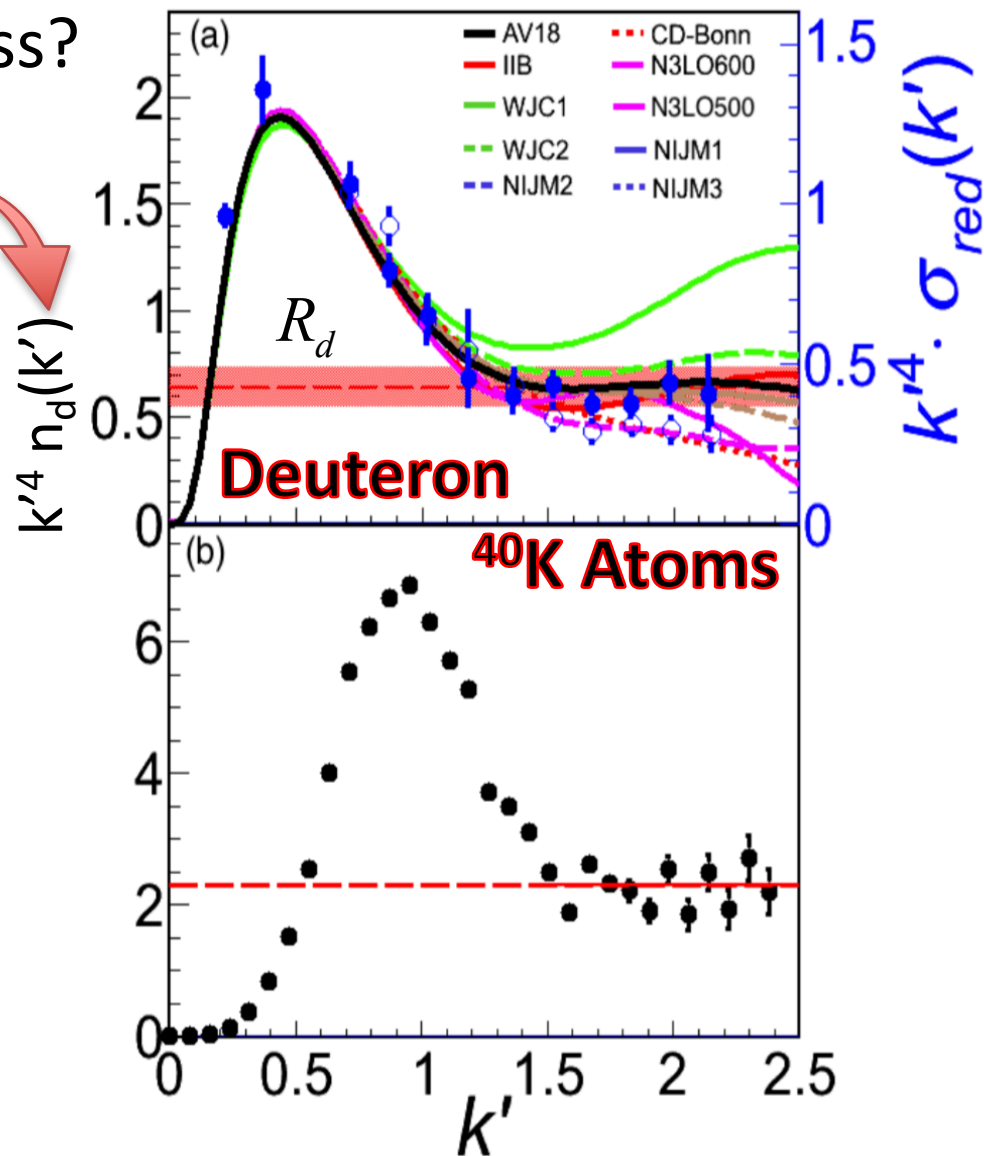
But Experiment Says.... Yes!



Is there $1/k^4$ scaling regardless?

$$1.5k_F < k < 3k_F$$

$$n_A(k) = a_2(A/d) \cdot n_d(k)$$



O. Hen et al. Phys. Rev. C **92**, 045205 (2015)



But Experiment Says.... Yes! (?)



Is there $1/k^4$ scaling regardless?

$$1.5k_F < k < 3k_F$$

$$n_A(k) = a_2(A/d) \cdot n_d(k)$$

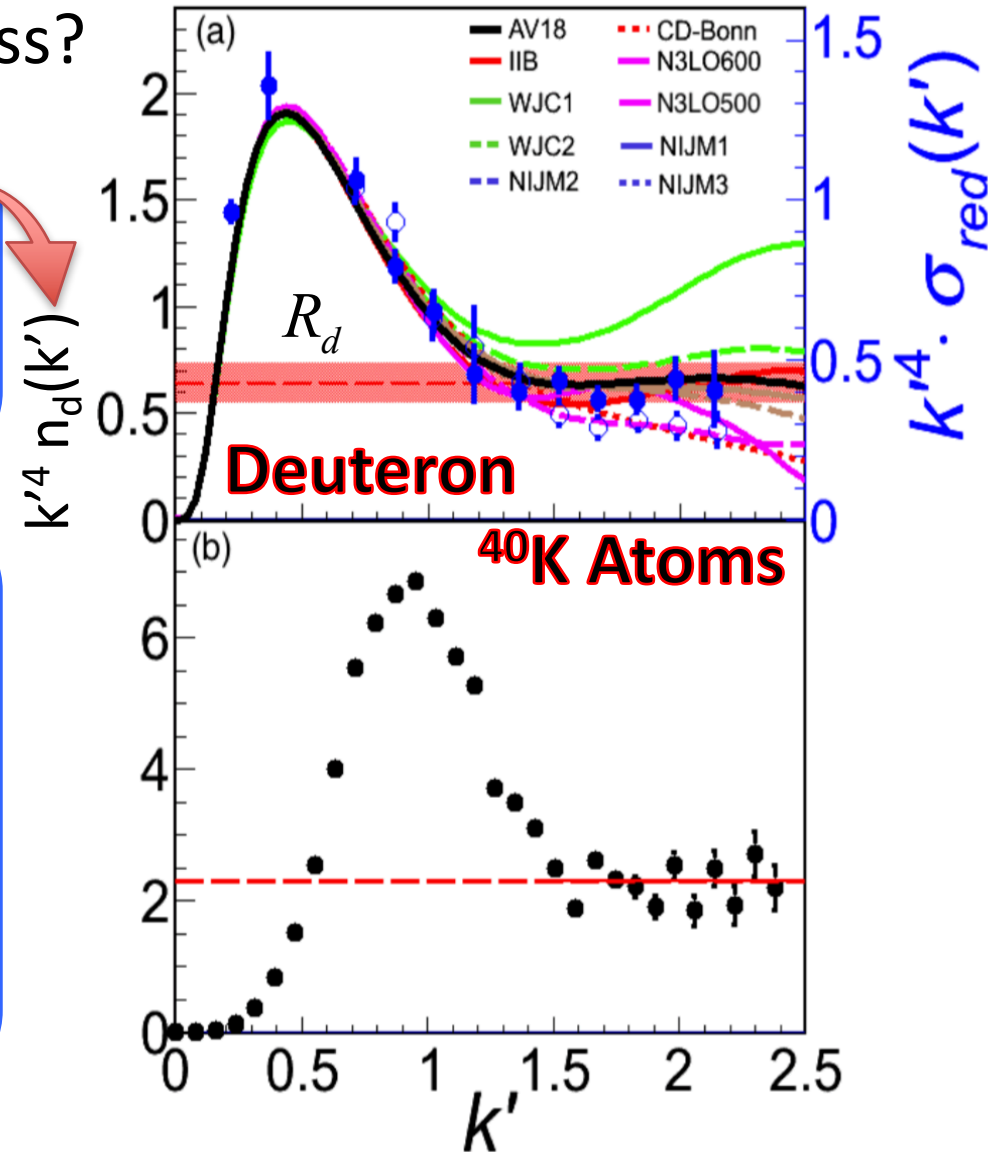
Why $1/k^4$?

Effect of the one pion exchange (OPE) contribution to the tensor potential acting in second order

$$(-B - H_0)|\Psi_D\rangle = V_T|\Psi_S\rangle$$

$$V_{00} = V_T(-B - H_0)^{-1} V_T$$

O. Hen et al. Phys. Rev. C **92**, 045205 (2015)





But Experiment Says.... Yes! (?)



Is there $1/k^4$ scaling regardless?

$$1.5k_F < k < 3k_F$$

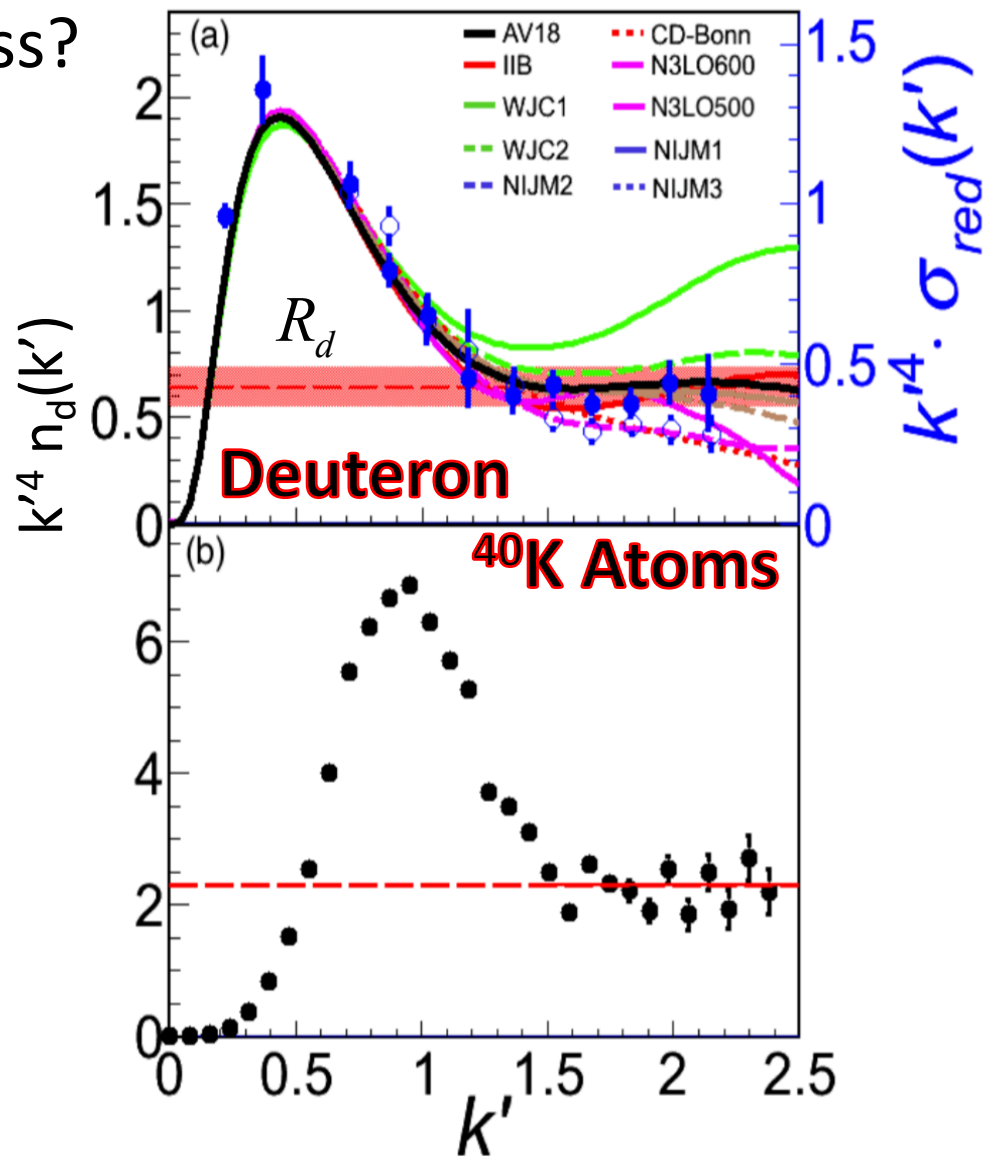
$$n_A(k) = a_2(A/d) \cdot n_d(k)$$



$$\frac{C}{k_F \cdot A} = a_2(A) \cdot R_d$$

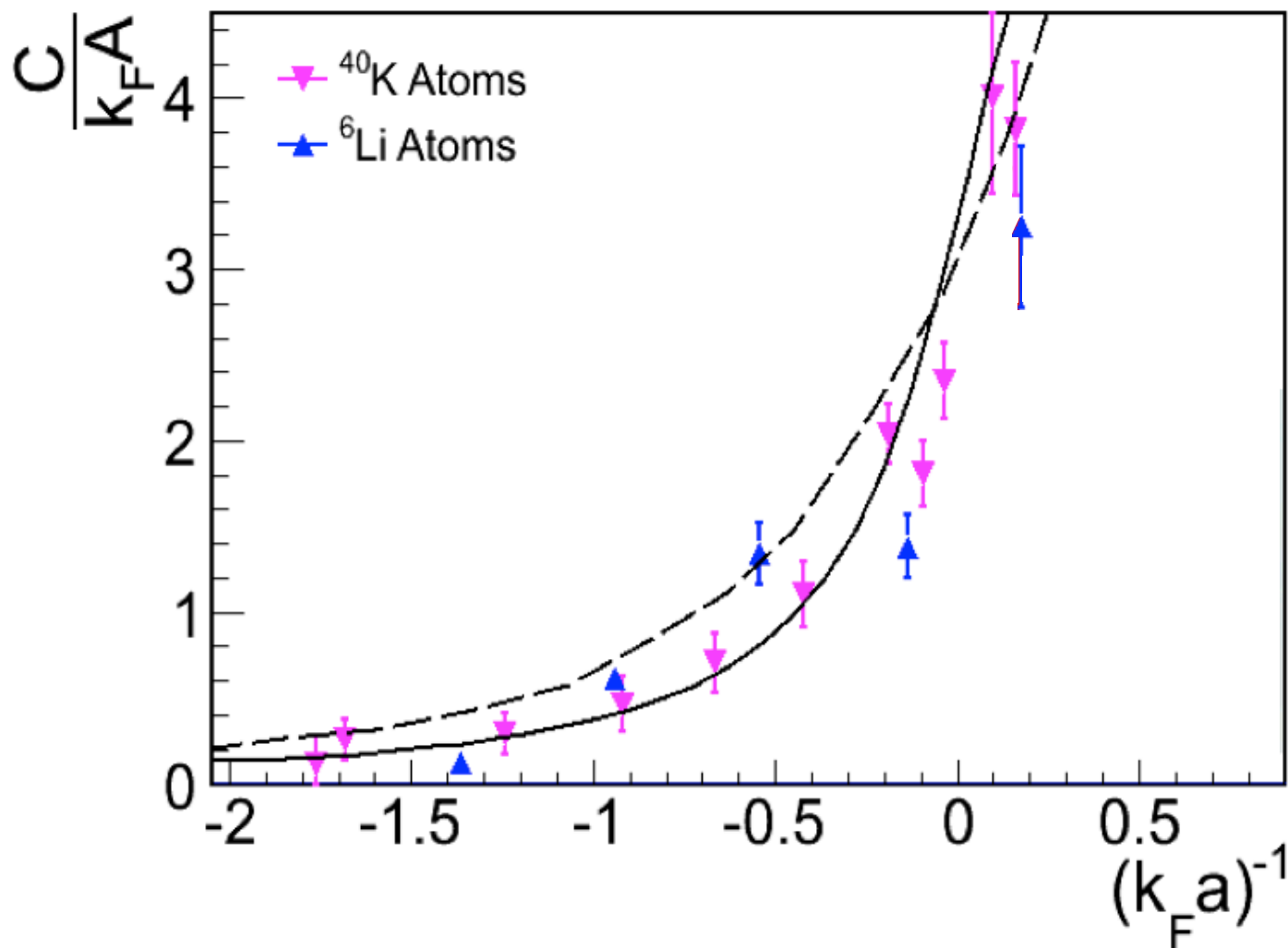
Nucleus	$a_2(A)$	$\frac{C}{k_F A}$
^{12}C	4.75 ± 0.16	3.04 ± 0.49
^{56}Fe	5.21 ± 0.20	3.33 ± 0.54
^{197}Au	5.16 ± 0.22	3.30 ± 0.53

O. Hen et al. Phys. Rev. C **92**, 045205 (2015)





Comparing with atomic systems



Stewart et al. Phys. Rev. Lett. **104**, 235301 (2010)

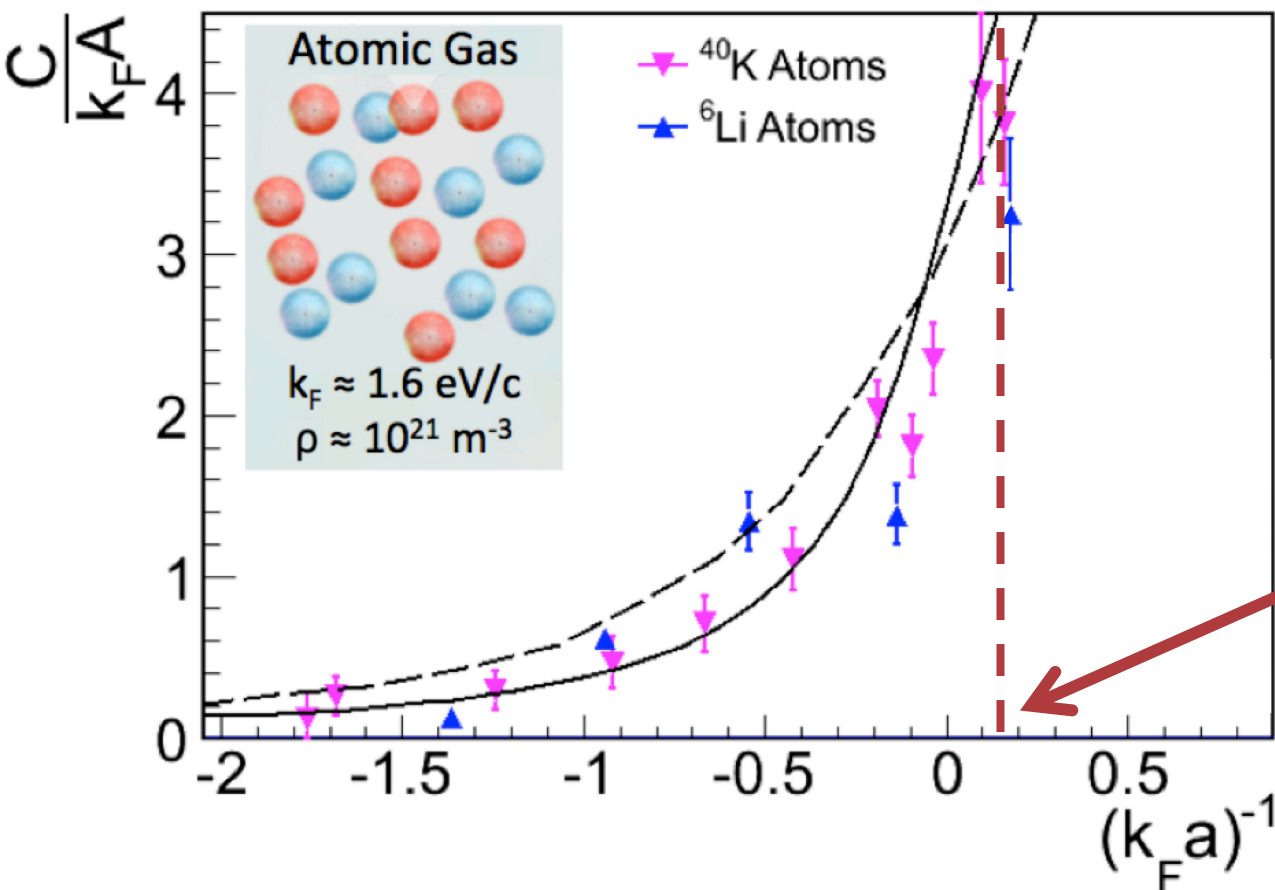
Kuhnle et al. Phys. Rev. Lett. **105**, 070402 (2010)



Comparing with atomic systems



Finding the same *dimensionless* interaction strength



For Nuclei:

$$k_F \approx 1.27 \text{ fm}^{-1}$$

$$a \approx 5.4 \text{ fm}$$

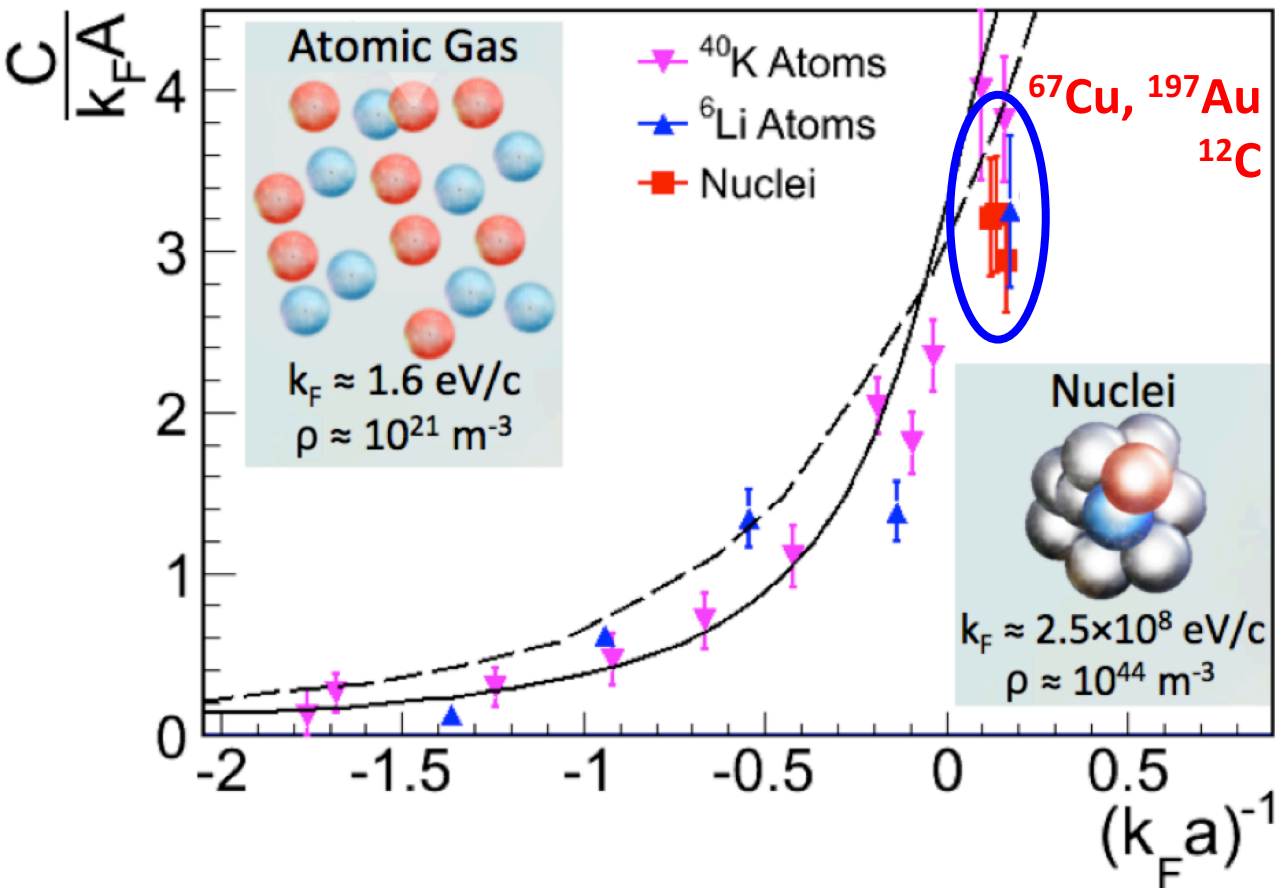
$$(k_F a)^{-1} \approx 0.15$$

Stewart et al. Phys. Rev. Lett. **104**, 235301 (2010)
Kuhnle et al. Phys. Rev. Lett. **105**, 070402 (2010)



Comparing with atomic systems

Equal contacts for equal interactions strength!



For Nuclei:

$$k_F \approx 1.27 \text{ fm}^{-1}$$

$$a \approx 5.4 \text{ fm}$$

$$(k_F a)^{-1} \approx 0.15$$

Nucleus	$\frac{C}{k_F A}$
^{12}C	3.04 ± 0.49
^{56}Fe	3.33 ± 0.54
^{197}Au	3.30 ± 0.53

$$\frac{C}{k_F \cdot A} = a_2(A) \cdot R_d$$

O. Hen et al. Phys. Rev. C **92**, 045205 (2015)
 Stewart et al. Phys. Rev. Lett. **104**, 235301 (2010)
 Kuhnle et al. Phys. Rev. Lett. **105**, 070402 (2010)



How can we reconcile the experimental observation with theory expectation?



i.e. is there a region in which the nuclear wave function fully factorizes?



Going Back to the Theory...

1. Generalize the contact formalism to nuclear systems.
2. Use it to make specific predictions of nuclear properties.
3. Check using experimental data and full many-body calculations.



The Contact and Universal Relations



Issue: Scale separation does not necessarily work in nuclear systems.

Solution: assume a more general form for the 2-body wavefunction.

Atomic System:

$$\Psi \xrightarrow{r_{ij} \rightarrow 0} (1/r_{ij} - 1/a) A_{ij}(\mathbf{R}_{ij}, \{\mathbf{r}_k\}_{k \neq i,j})$$

Nuclear System:

$$\Psi \xrightarrow{r_{ij} \rightarrow 0} (\quad \boldsymbol{\varphi}(r)_{ij} \quad) A_{ij}(\mathbf{R}_{ij}, \{\mathbf{r}_k\}_{k \neq i,j})$$



The Contact and Universal Relations



Issue: Scale separation does not necessarily work in nuclear systems.

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Atomic System:

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Nuclear System:

$$\Psi \xrightarrow{r_{ij} \rightarrow 0} (\quad \boldsymbol{\varphi}(r)_{ij} \quad) A_{ij}(\mathbf{R}_{ij}, \{\mathbf{r}_k\}_{k \neq i,j})$$



known solution for the two-body (nuclear) problem



Factorization in Nuclei

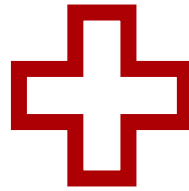


Consider the factorized wave function:

$$\Psi \xrightarrow{r_{ij} \rightarrow 0} \sum_{\alpha} \varphi_{\alpha}(\mathbf{r}_{ij}) A_{ij}^{\alpha}(\mathbf{R}_{ij}, \{\mathbf{r}_k\}_{k \neq i,j})$$

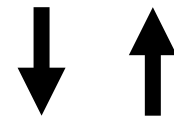
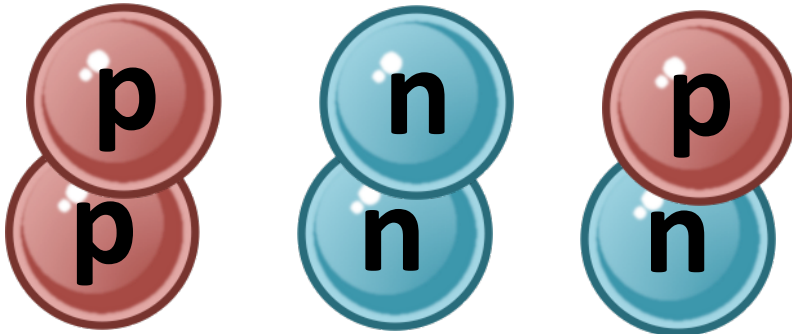
In nuclear physics we have 3 possible types of pairs:

$$ij = \{pp, nn, pn\}$$



For each pair we have different channels

$$\alpha = (s,l)jm$$





Factorization in Nuclei



Consider the factorized wave function:

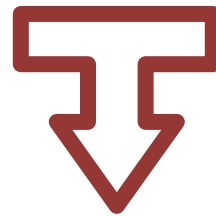
$$\Psi \xrightarrow{r_{ij} \rightarrow 0} \sum_{\alpha} \varphi_{\alpha}(\mathbf{r}_{ij}) A_{ij}^{\alpha}(\mathbf{R}_{ij}, \{\mathbf{r}_k\}_{k \neq i,j})$$

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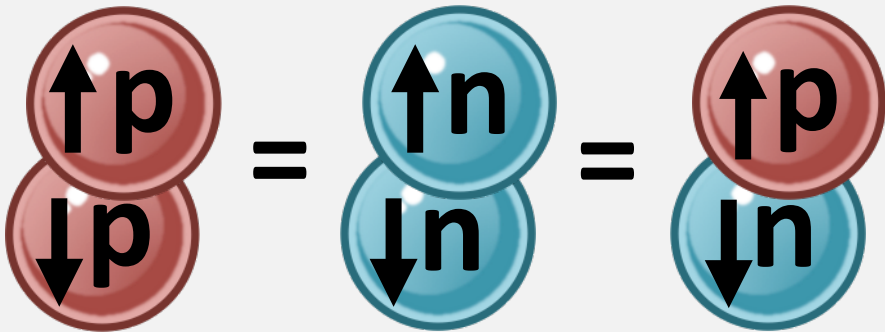
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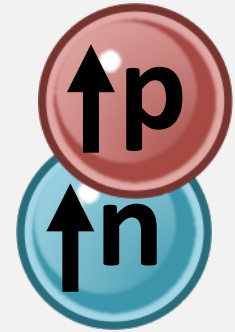


Reduced to 2 contacts by imposing $L=0$ and symmetry considerations

$S=0$



$S=1$





Relating to Momentum Space



$$\Psi \xrightarrow{r_{ij} \rightarrow 0} \sum_{\alpha} \varphi_{\alpha}(\mathbf{r}_{ij}) A_{ij}^{\alpha}(\mathbf{R}_{ij}, \{\mathbf{r}_k\}_{k \neq i,j})$$



One Body:

$$n_p(\mathbf{k}) = \sum_{\alpha} |\tilde{\varphi}_{pp}^{\alpha}(\mathbf{k})|^2 2\mathbf{C}_{pp}^{\alpha} + \sum_{\alpha} |\tilde{\varphi}_{pn}^{\alpha}(\mathbf{k})|^2 \mathbf{C}_{pn}^{\alpha}$$



2-Body momentum distributions



- One Body momentum distribution $[n_N(k)]$:
Probability to find a nucleon, N , in the nucleus with momentum k .
- Two Body momentum distribution $[n_{NN}(q,Q)]$:
Probability to find a NN pair in the nucleus with relative (c.m.) momentum q (Q).

$n_{NN}(q,Q)$ – computational Frontier!



Momentum Space Factorization



$$\Psi \xrightarrow{r_{ij} \rightarrow 0} \sum_{\alpha} \varphi_{\alpha}(\mathbf{r}_{ij}) A_{ij}^{\alpha}(\mathbf{R}_{ij}, \{\mathbf{r}_k\}_{k \neq i,j})$$



One Body:

$$n_p(\mathbf{k}) = \sum_{\alpha} |\tilde{\varphi}_{pp}^{\alpha}(\mathbf{k})|^2 2\mathbf{C}_{pp}^{\alpha} + \sum_{\alpha} |\tilde{\varphi}_{pn}^{\alpha}(\mathbf{k})|^2 \mathbf{C}_{pn}^{\alpha}$$

Two body:

$$F_{ij}(\mathbf{k}) = \sum_{\alpha} |\tilde{\varphi}_{ij}^{\alpha}(\mathbf{k})|^2 \mathbf{C}_{ij}^{\alpha}$$

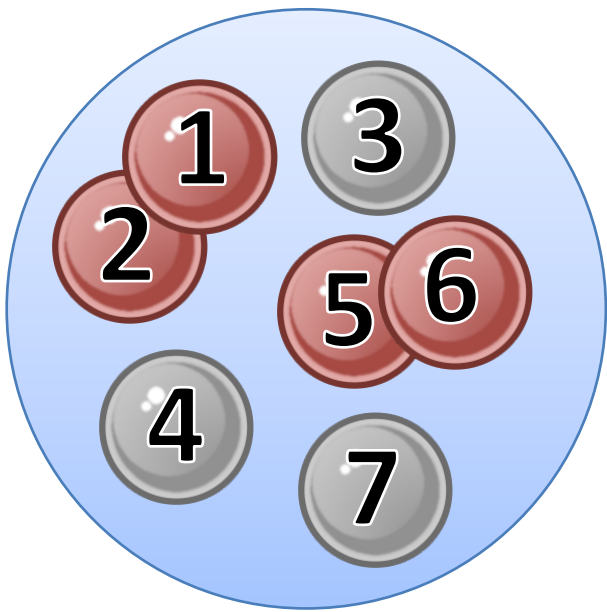
Clearly:

$$n_p(\mathbf{k}) \xrightarrow{\mathbf{k} \rightarrow \infty} 2F_{pp}(\mathbf{k}) + F_{pn}(\mathbf{k})$$



Two-Body Momentum Distributions

- $n_{NN}(q,Q)$ – Mathematical object that counts all possible NN pairs, regardless of their state:



Consider all NN pairs:

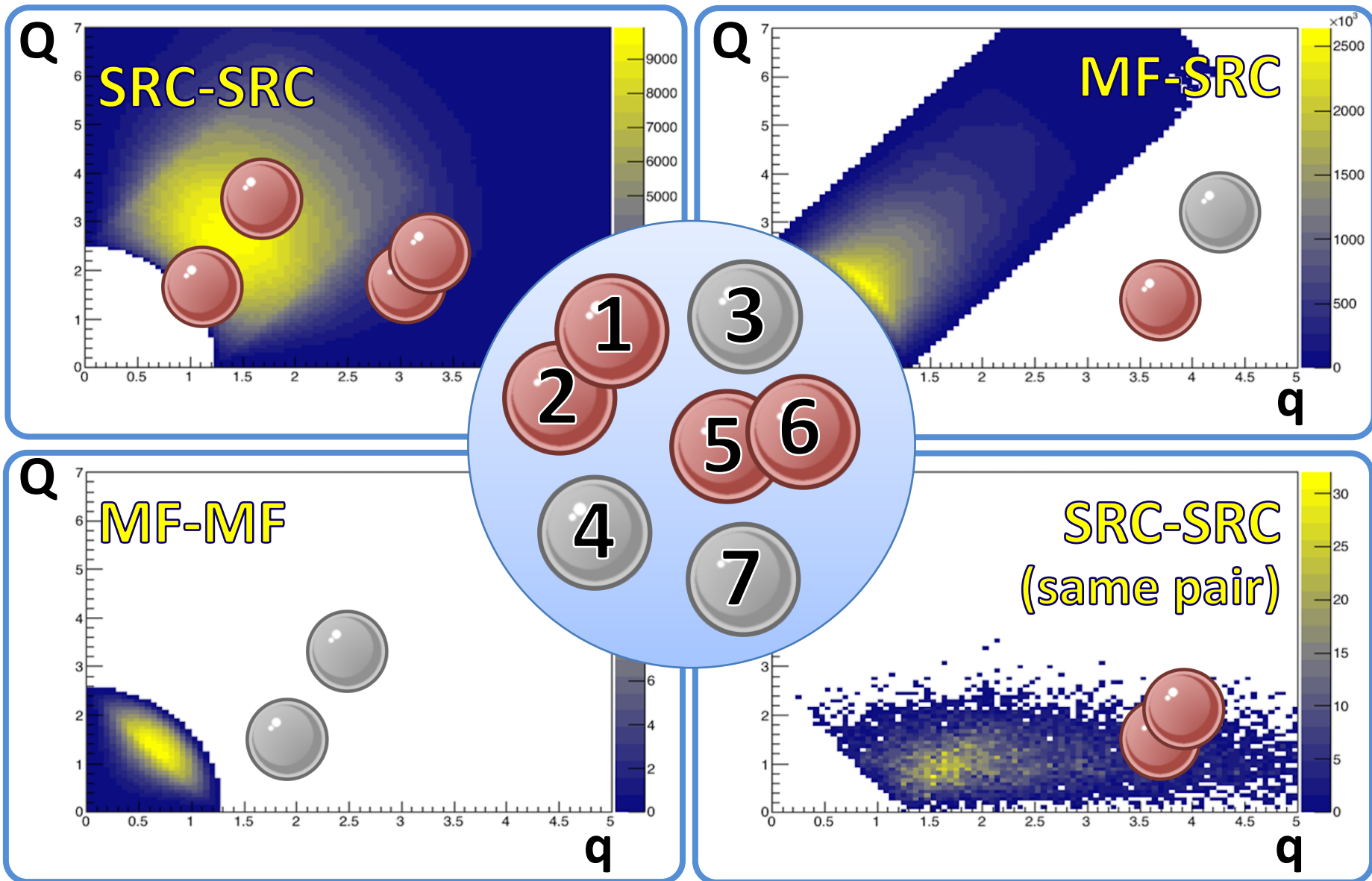
<u>1-2</u>	2-3	3-4	4-5	<u>5-6</u>	6-7
1-3	2-4	3-5	4-6	5-7	
1-4	<u>2-5</u>	3-6	4-7		
<u>1-5</u>	<u>2-6</u>	3-7			
<u>1-6</u>	2-7				
1-7					



$n_{NN}(q,Q)$

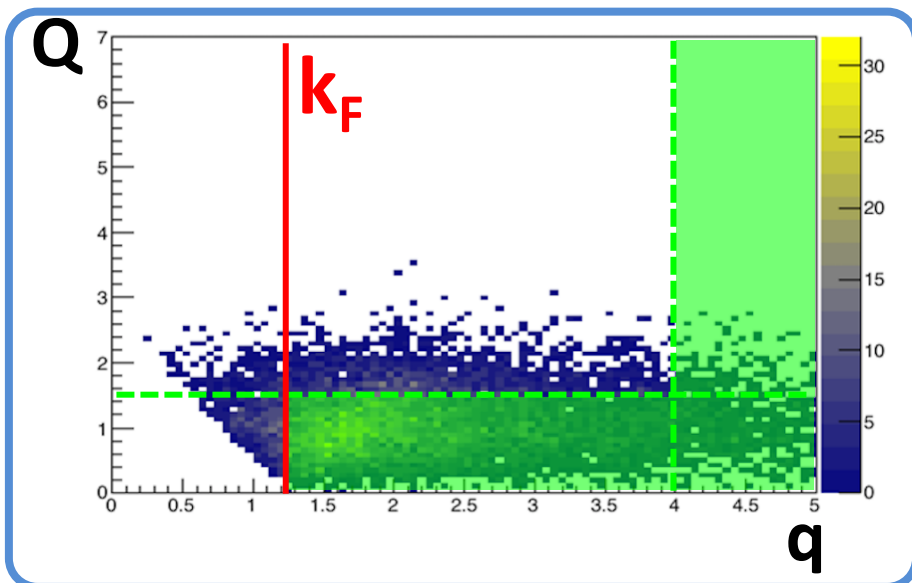
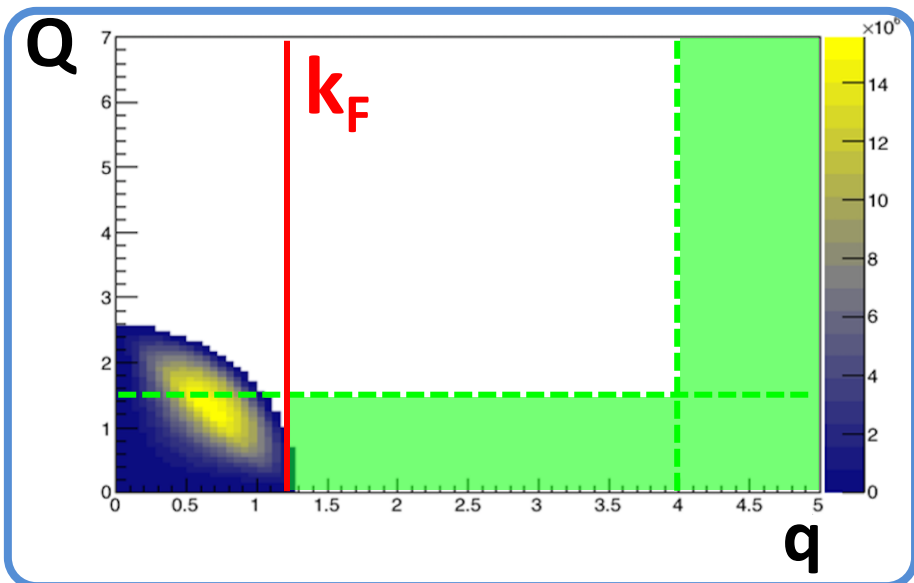
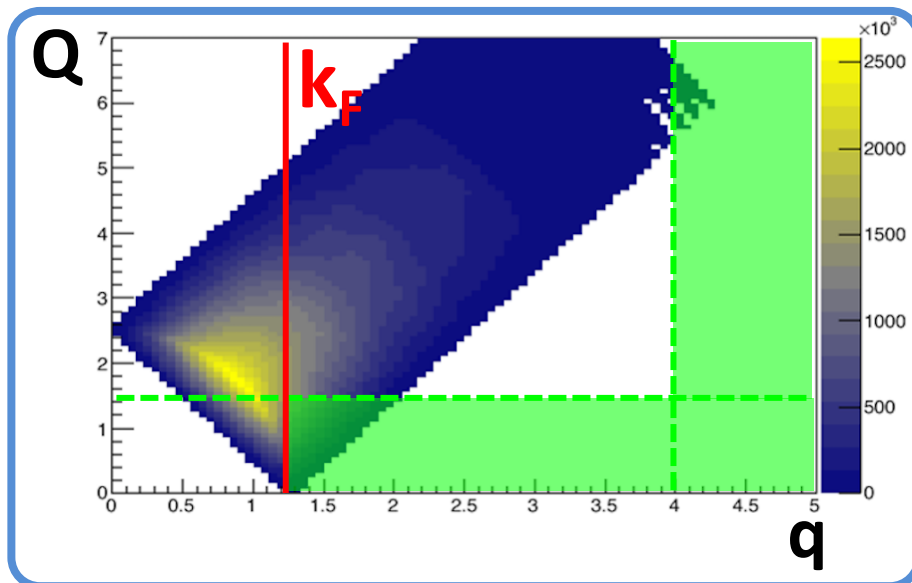
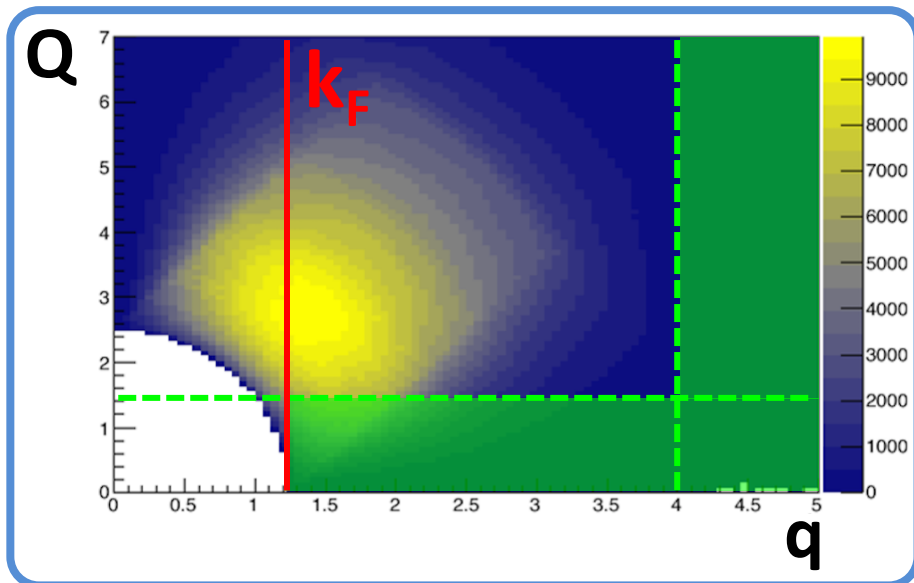


Toy model to the rescue





Toy model to the rescue

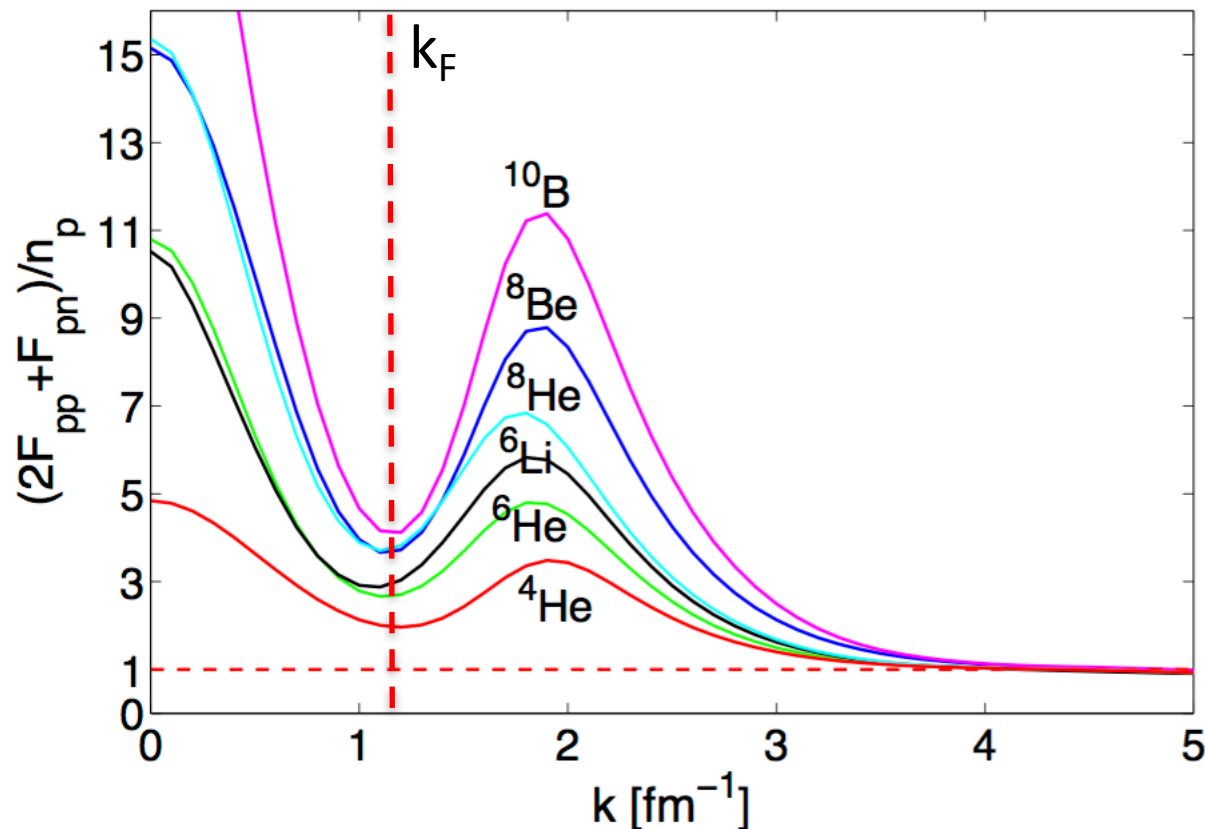




Two-Body Scaling for High q

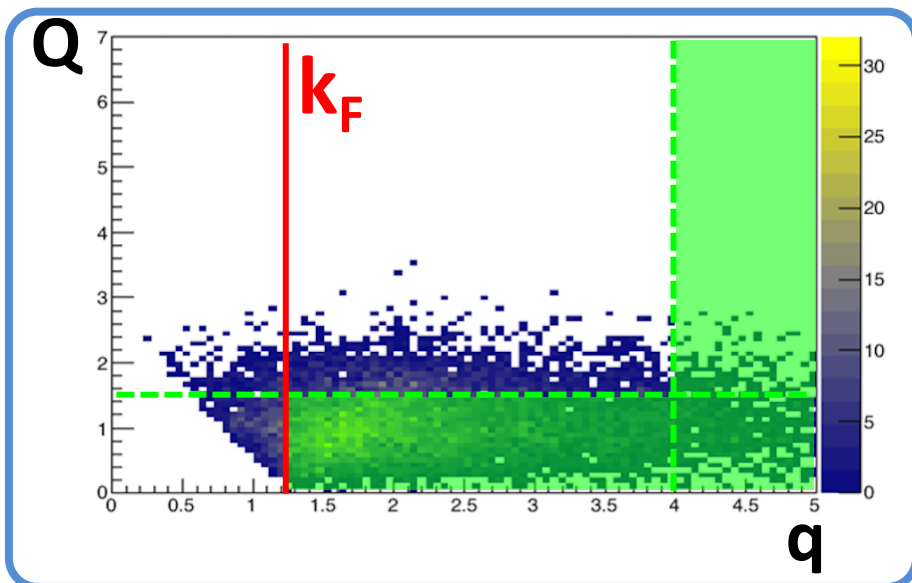
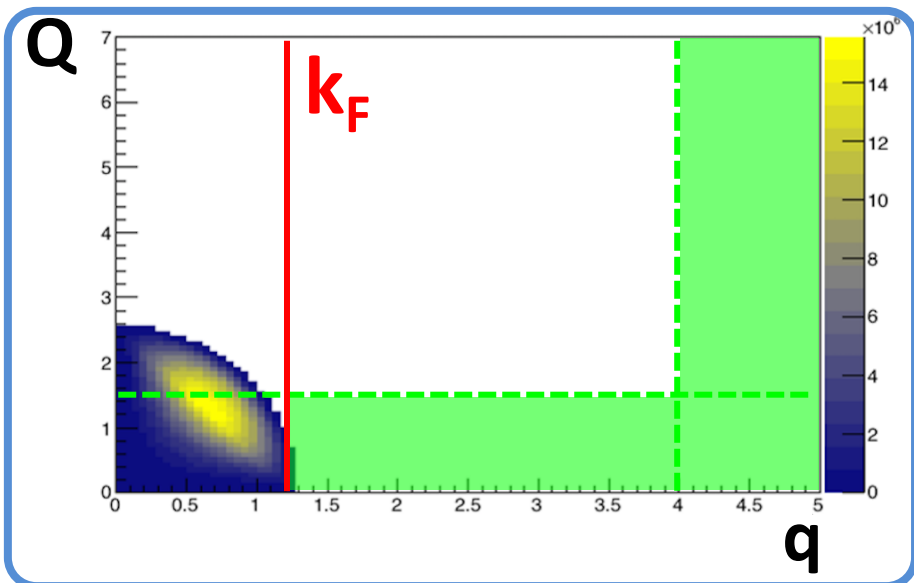
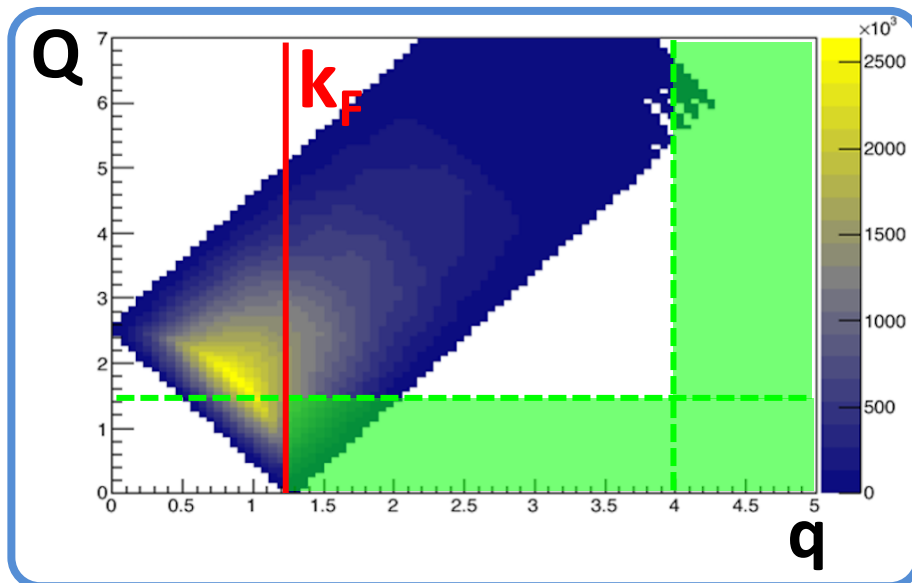
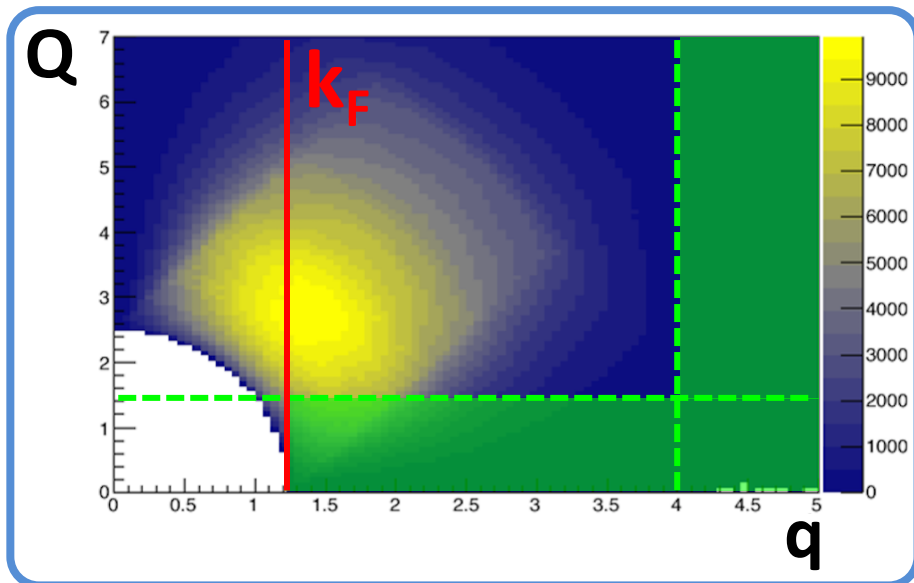


- Weiss and Barnea (PRC 2015): contact interactions dominate when $n_{pn}(q) + 2n_{pp}(q) = n_p(k)$





Toy model to the rescue

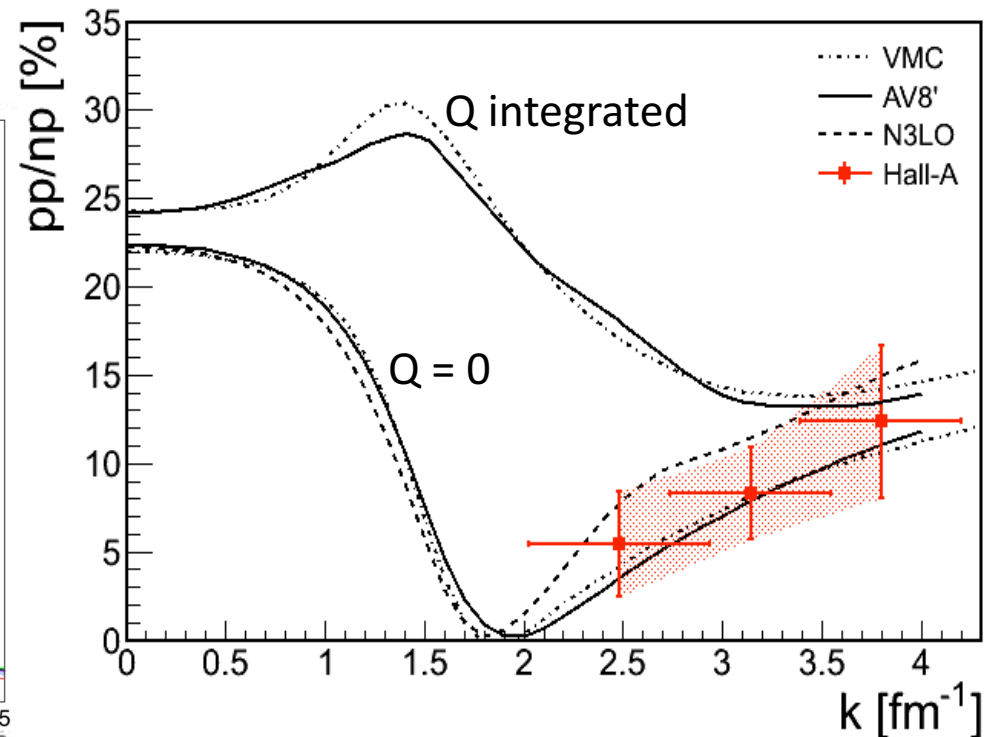
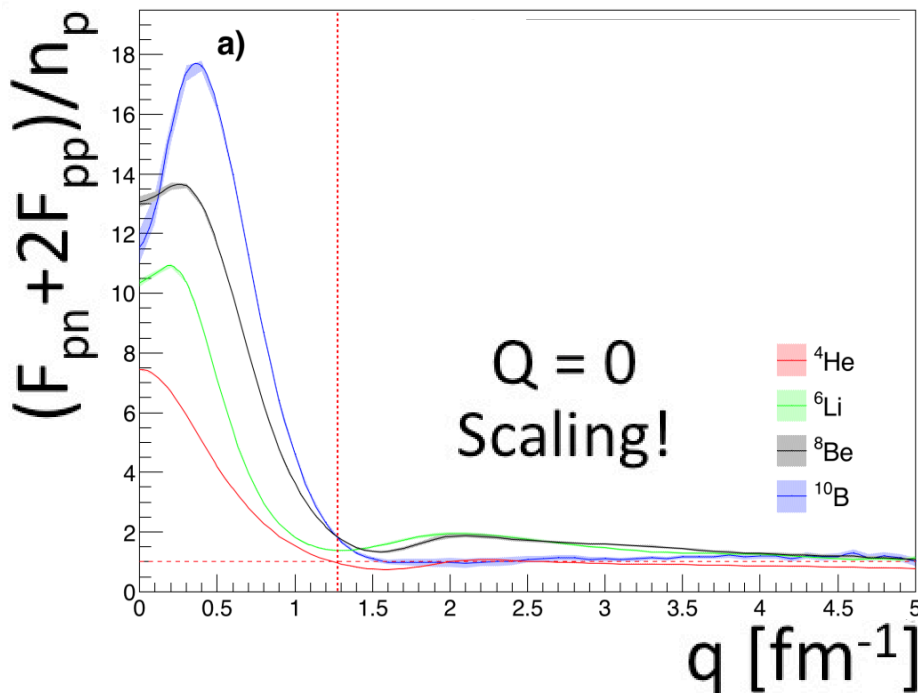




Two-Body Scaling for Q=0



- Restricting $Q=0$ restores scaling starting from $k > k_F$ AND gives consistent results with experimental data!



SRC pairs are consistent with $Q = 0$ *back-to-back* pairs

R. Wiringa et al., Phys. Rev. C 89, 024305 (2014).

T. Neff, H. Feldmeier and W. Horiuchi, Phys. Rev. C 92, 024003 (2015).

I. Korover, N. Muangma, and O. Hen et al., Phys. Rev. Lett 113, 022501 (2014).

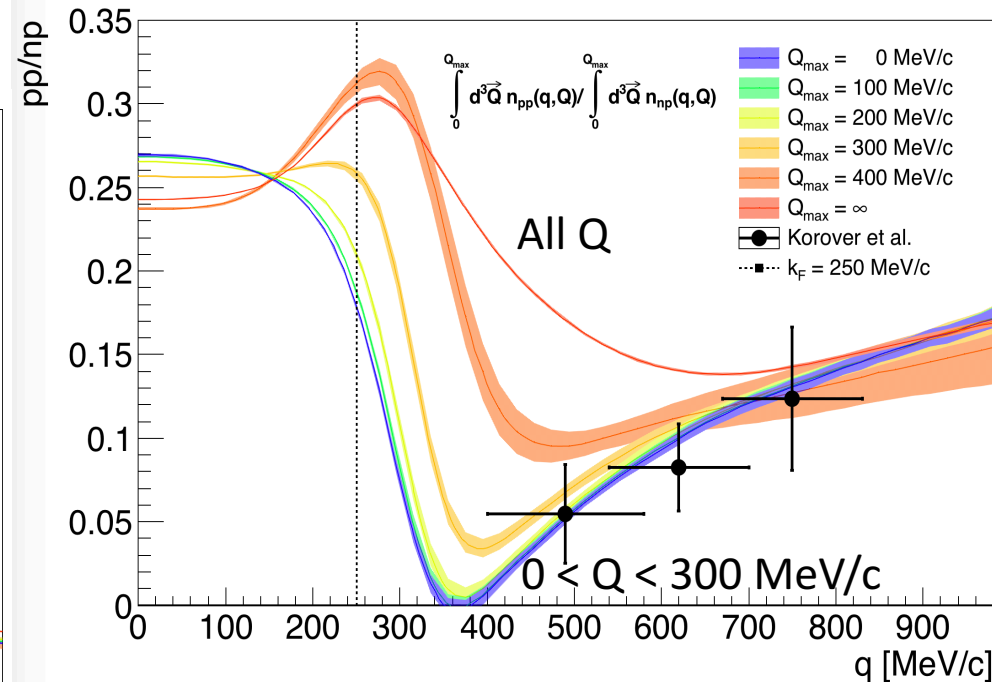
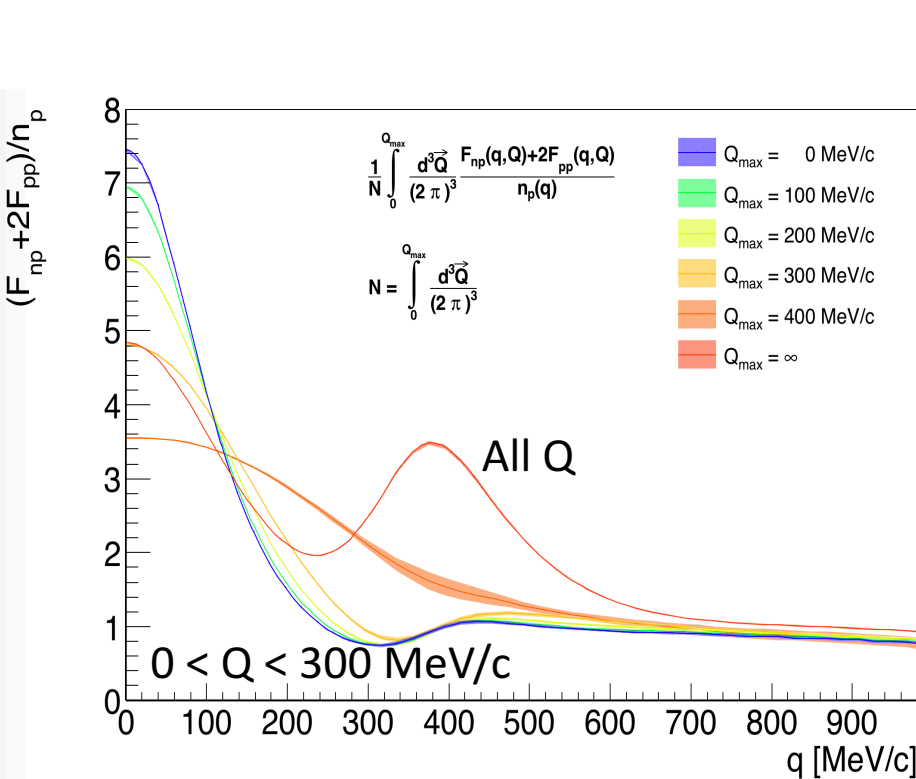
Weiss, Cruz-Torres, Barnea, Piasetzky and Hen, arXiv 1612.00923 (2016)



Two-Body Scaling for Low Q



- Restricting $Q=0$ restores scaling starting from $k > k_F$ AND gives consistent results with experimental data!



SRC pairs are consistent with $Q \leq k_F$ *back-to-back* pairs

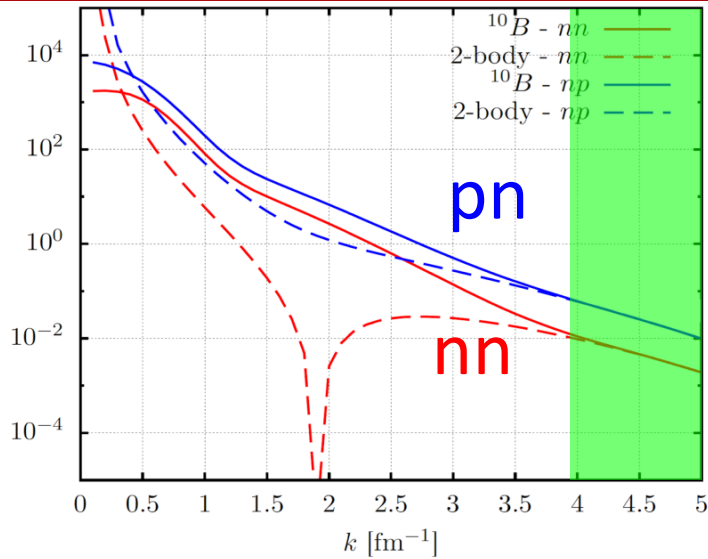
- R. Wiringa et al., Phys. Rev. C 89, 024305 (2014).
 T. Neff, H. Feldmeier and W. Horiuchi, Phys. Rev. C 92, 024003 (2015).
 I. Korover, N. Muangma, and O. Hen et al., Phys. Rev. Lett 113, 022501 (2014).
 Weiss, Cruz-Torres, Barnea, Piasetzky and Hen, arXiv 1612.00923 (2016)

Extracting the nuclear contact(s)





Extracting the Contacts

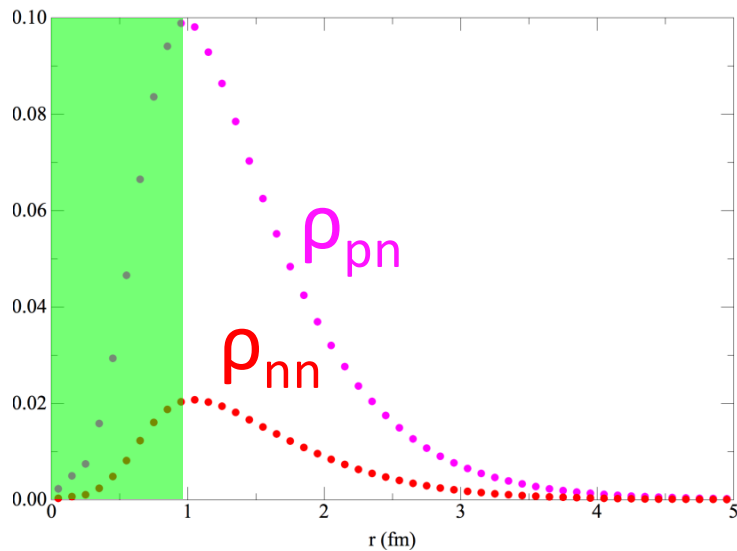


2-Body momentum distributions

$$F_{pn}(\mathbf{k}) \xrightarrow{k \rightarrow \infty} |\varphi_{pn}^0(\mathbf{k})|^2 C_{pn}^0 + |\varphi_{pn}^d(\mathbf{k})|^2 C_{pn}^d$$

$$F_{nn}(\mathbf{k}) \xrightarrow{k \rightarrow \infty} |\varphi_{nn}^0(\mathbf{k})|^2 C_{nn}^0$$

Fitting range $\sim 4\text{-}5 \text{ fm}^{-1}$



2-Body density distributions

$$\rho_{pn}(\mathbf{r}) \xrightarrow{r \rightarrow 0} |\varphi_{pn}^0(\mathbf{r})|^2 C_{pn}^0 + |\varphi_{pn}^d(\mathbf{r})|^2 C_{pn}^d$$

$$\rho_{nn}(\mathbf{r}) \xrightarrow{r \rightarrow 0} |\varphi_{nn}^0(\mathbf{r})|^2 C_{nn}^0$$

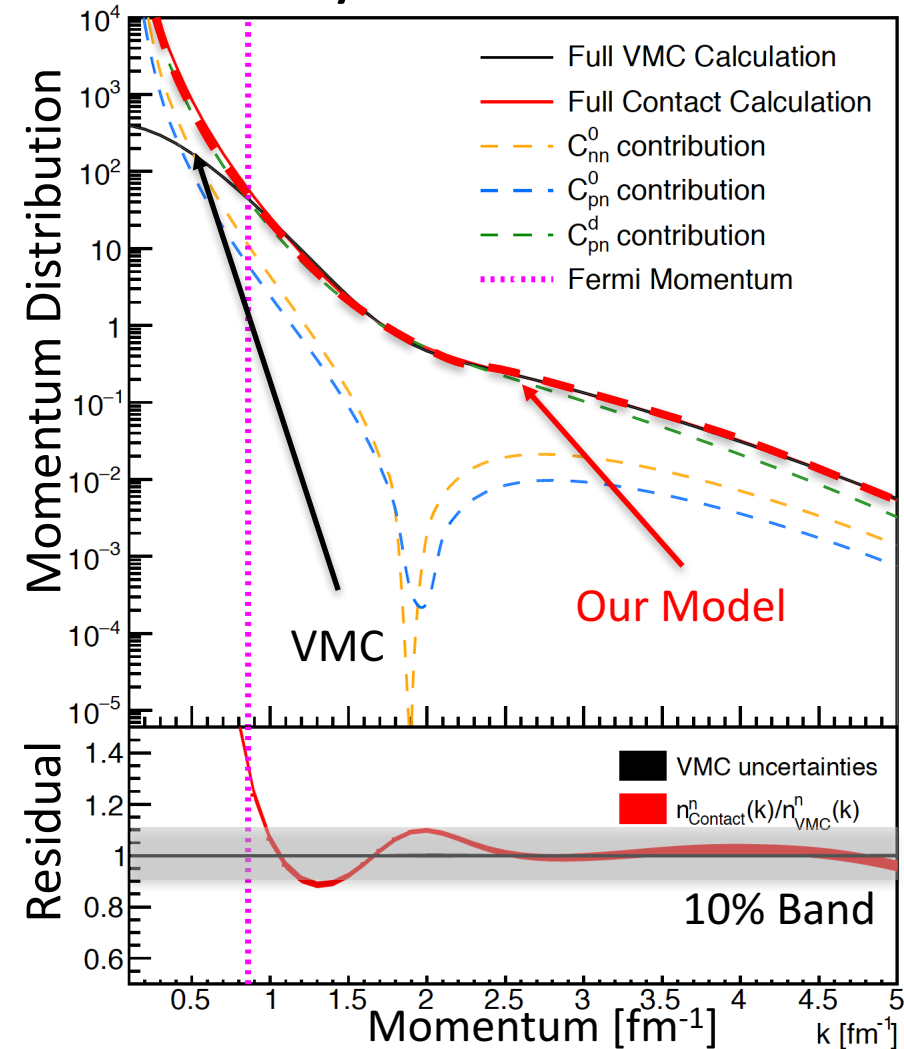
Fitting range $\sim 0.25\text{-}1 \text{ fm}$



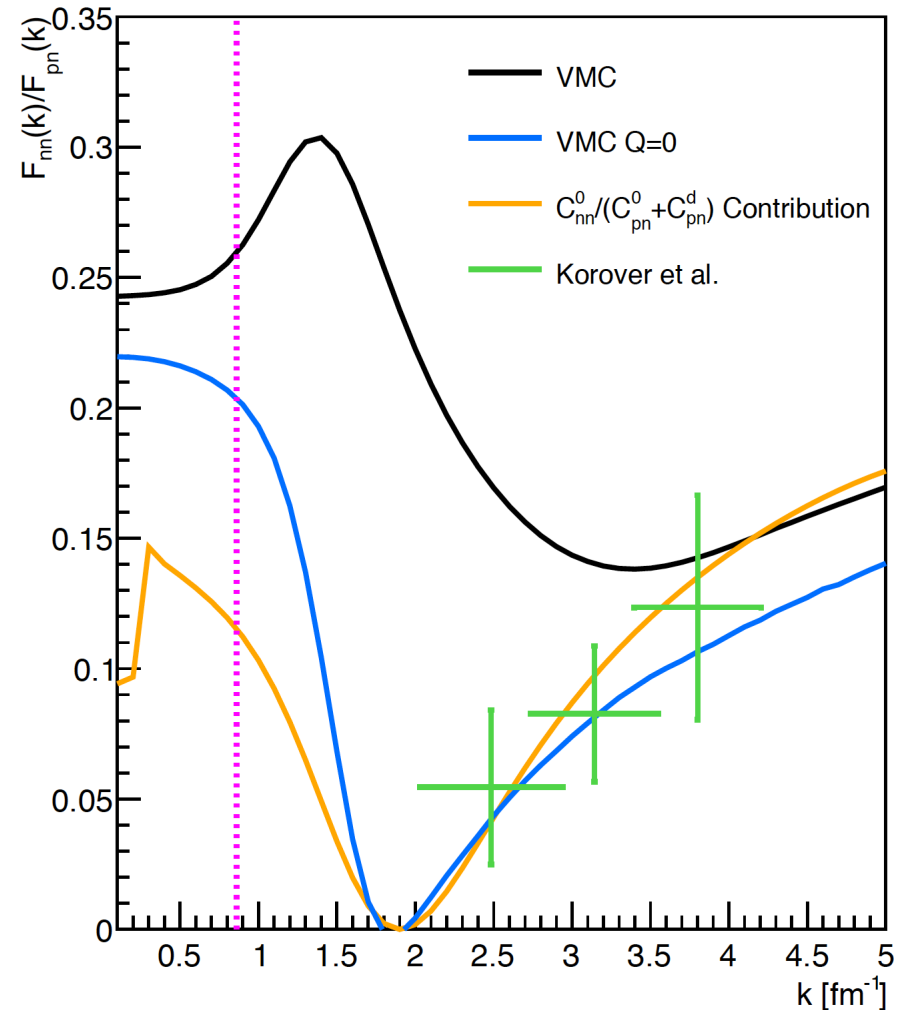
^4He Results



1-body Momentum dist.



pp / np ratio



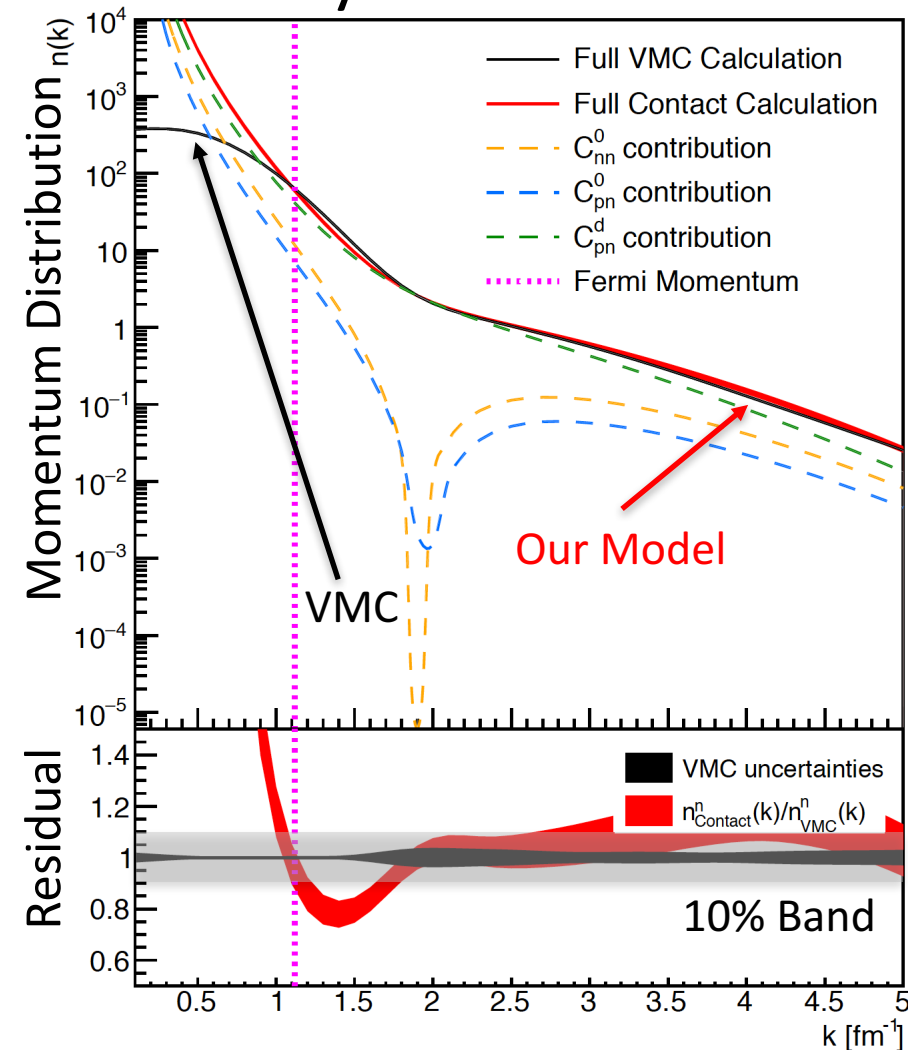
Weiss, Cruz-Torres, Barnea, Piasezky and Hen, arXiv 1612.00923 (2016)



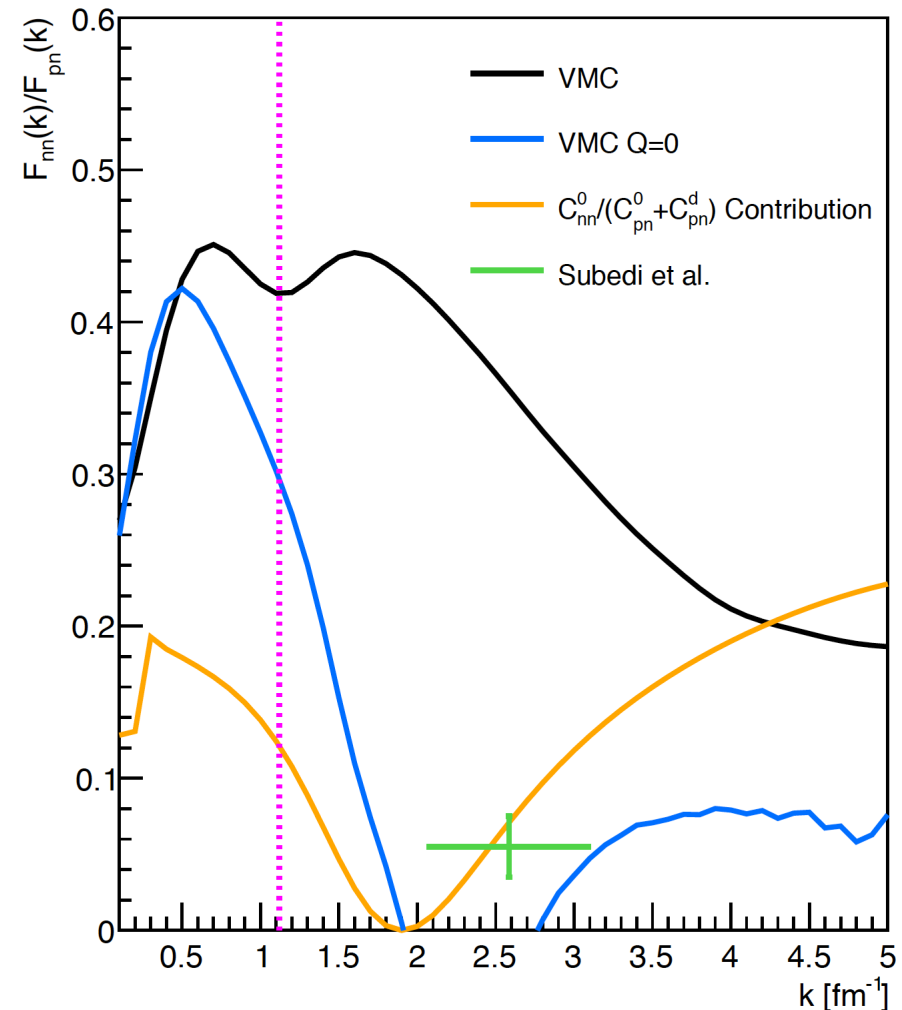
^{12}C Results



1-body Momentum dist.



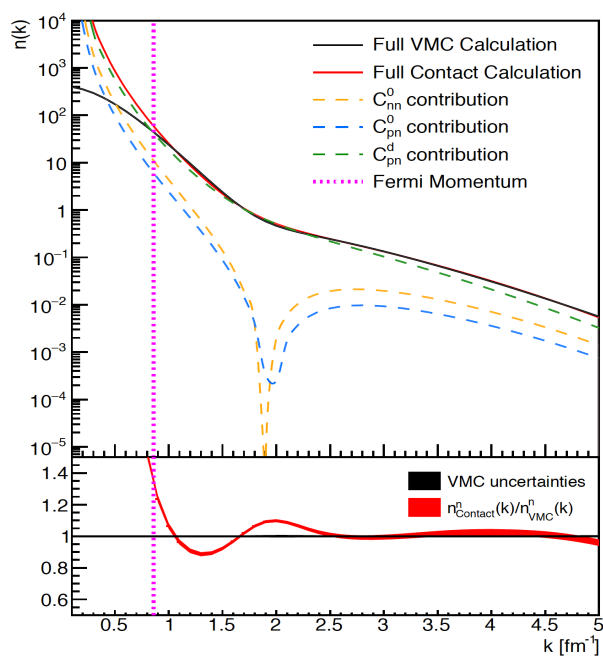
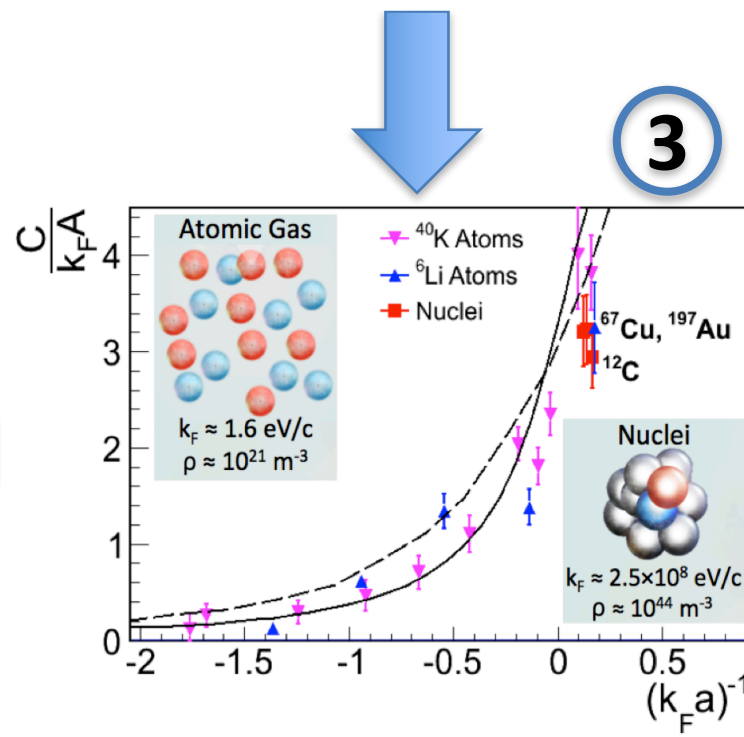
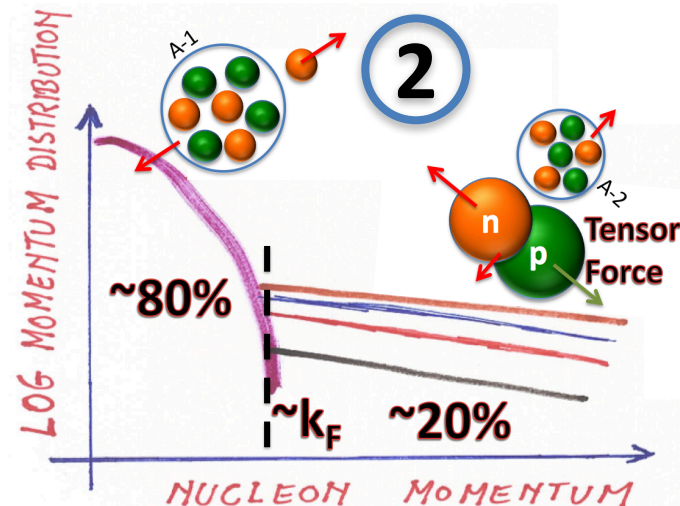
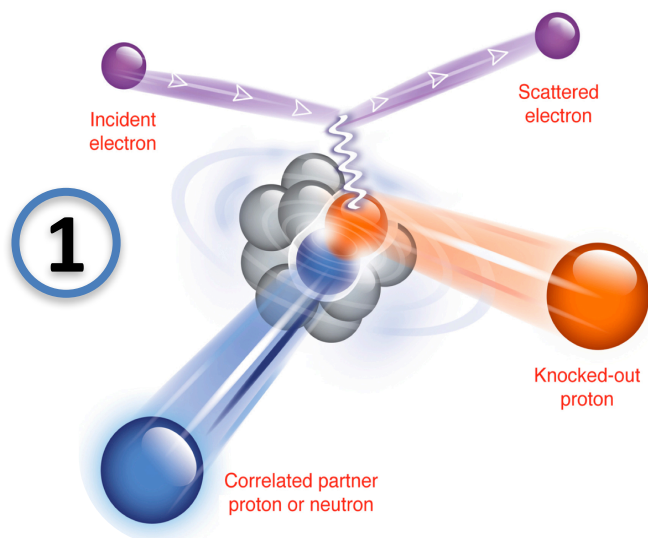
pp / np ratio



Weiss, Cruz-Torres, Barnea, Piasezky and Hen, arXiv 1612.00923 (2016)

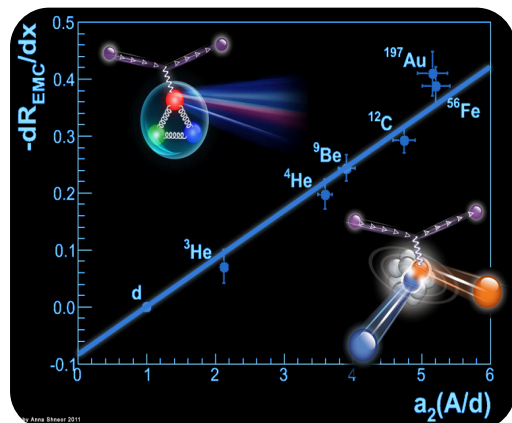


Summery: Exp. to Pheno. to Theory

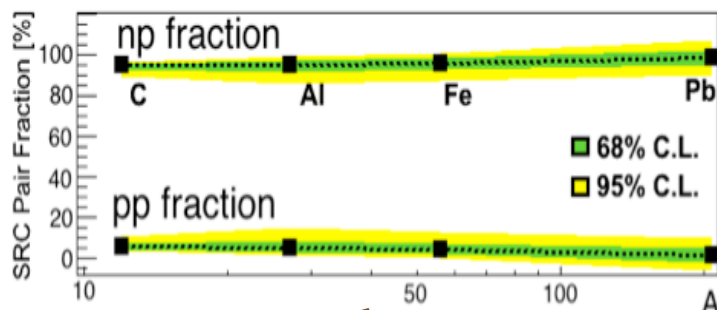




Why SRC?

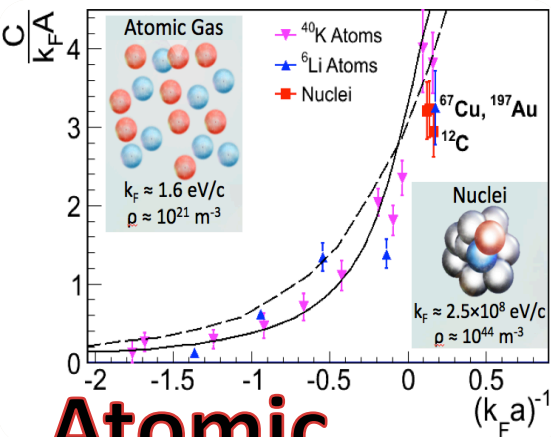
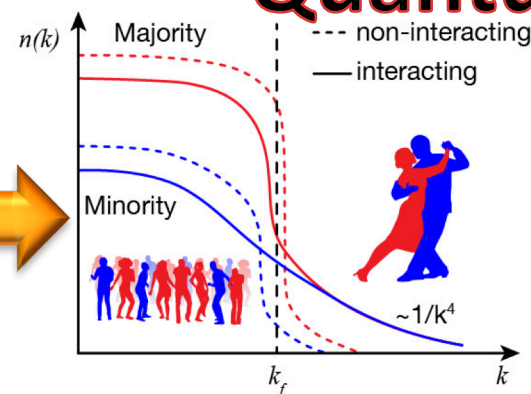


Particle



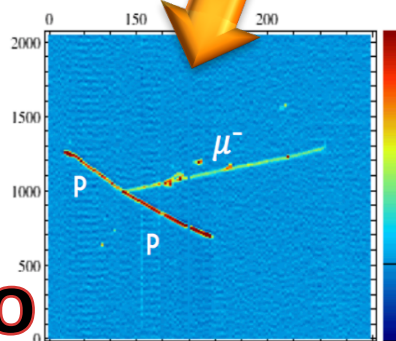
Nuclear

Quantum

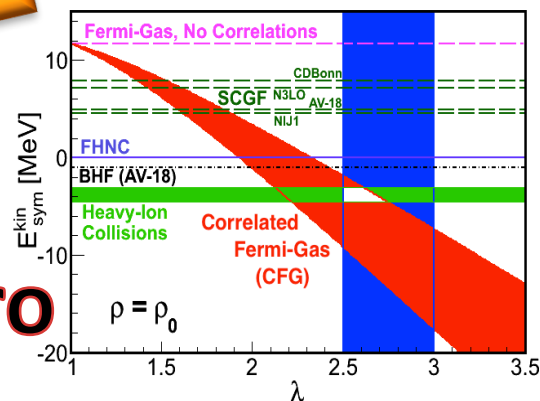


Atomic

Neutrino

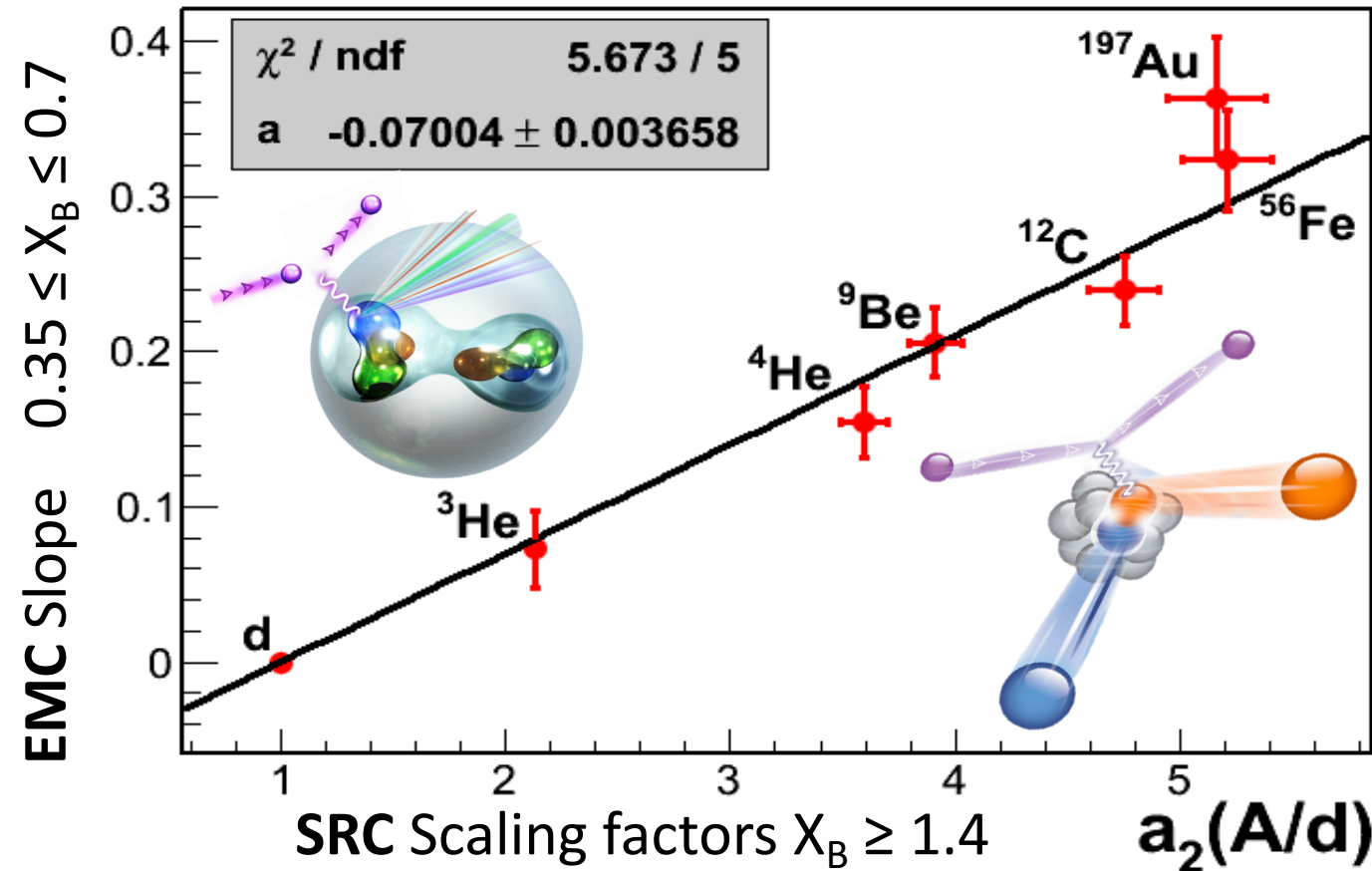


Astro

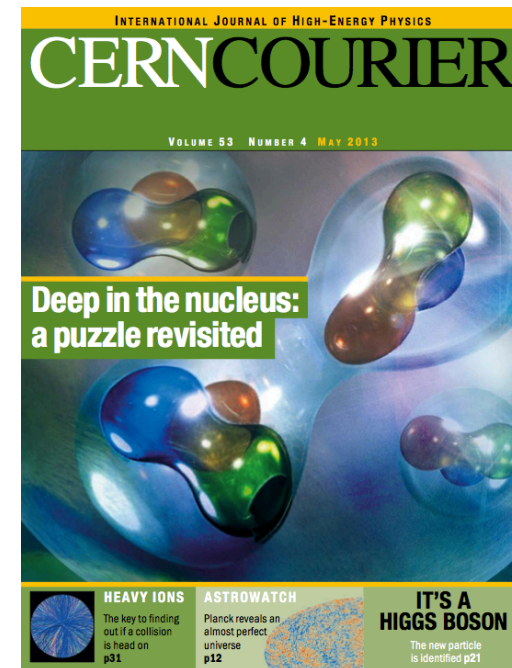




EMC-SRC Correlation



See G. Miller
talk yesterday



O. Hen et al., Int. J. Mod. Phys. E. **22**, 1330017 (2013).

O. Hen et al., Phys. Rev. C **85** (2012) 047301.

L. B. Weinstein, E. Piasetzky, D. W. Higinbotham, J. Gomez, O. Hen, R. Shneor, Phys. Rev. Lett. **106** (2011) 052301.





Nucleon-Nucleon Correlations and the Quarks Within

Or Hen

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Gerald A. Miller

*Department of Physics,
University of Washington, Seattle,
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Eli Piassetzky

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Lawrence B. Weinstein

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Old Dominion University, Norfolk,
VA 23529*

(Dated: November 2, 2016)

- conventional (non-quark) nuclear physics cannot account for the EMC effect
- models need to include nucleon modification to account for the EMC effect. These models can modify the structure of either:
 - mean field nucleons, or
 - nucleons belonging to SRC pairs.
- there is a phenomenological connection between the strength of the EMC effect and the probability that a nucleon belongs to a two-nucleon SRC pair ($a_2(A)$), see Fig. 33.
- the influence of SRC pairs can account for the EMC-SRC correlation because both effects are driven by high virtuality nucleons with $p^2 \neq M^2$,
- the connection between the EMC effect and the coefficients $a_2(A)$ has been derived using two completely different theories, so that this connection is no accident
- nuclei must contain a small percentage of baryons that are not nucleons. Such baryons exist in the short-ranged correlations and are the source of the EMC effect.

Short Range Correlations and the EMC Effect in Effective Field Theory

Jiunn-Wei Chen,^{1,2,*} William Detmold,^{2,†} Joel E. Lynn,^{3,4,‡} and Achim Schwenk^{3,4,5,§}

¹Department of Physics, CTS and LeCosPA, National Taiwan University, Taipei 10617, Taiwan

²Center for Theoretical Physics, Massachusetts Institute of Technology, Cambridge, MA 02139, USA

³Institut für Kernphysik, Technische Universität Darmstadt, 64289 Darmstadt, Germany

⁴ExtreMe Matter Institute EMMI, GSI Helmholtzzentrum für Schwerionenforschung GmbH, 64291 Darmstadt, Germany

⁵Max-Planck-Institut für Kernphysik, Saupfercheckweg 1, 69117 Heidelberg, Germany

Gerald A. Miller

count for the EMC effect
the structure of either

arXiv: 1607.03065 (2016)

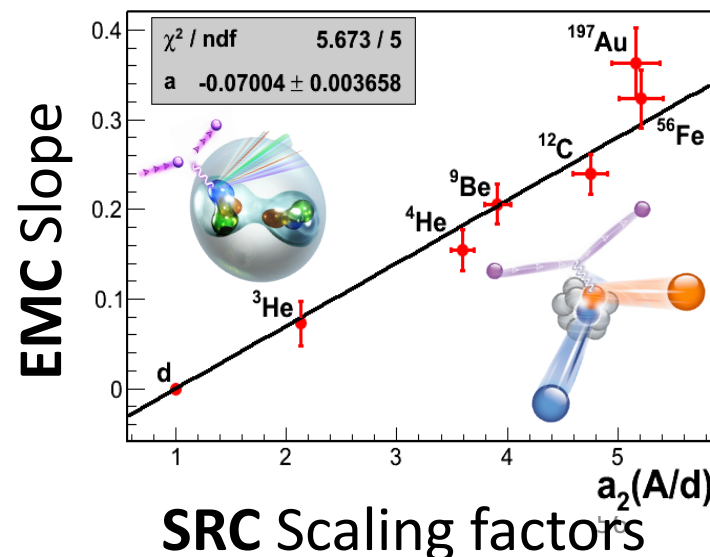
EFT description of bound nucleon structure:

$$F_2^A(x, Q^2)/A = F_2^N(x, Q^2) + g_2(A, \Lambda) f_2(x, Q^2, \Lambda).$$

$$g_2(A, \Lambda) = \frac{1}{A} \langle A | (N^\dagger N)^2 | A \rangle_\Lambda$$

The contact...

$$a_2(A, x > 1) = \frac{g_2(A, \Lambda)}{[\text{SRC Scaling Factor}] g_2(2, \Lambda)}$$





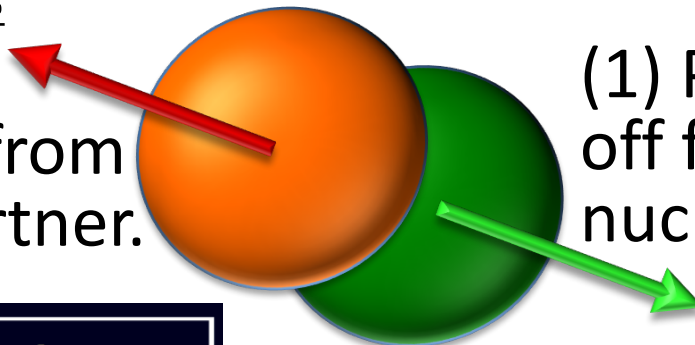
Tagged Structure Functions (JLab12)



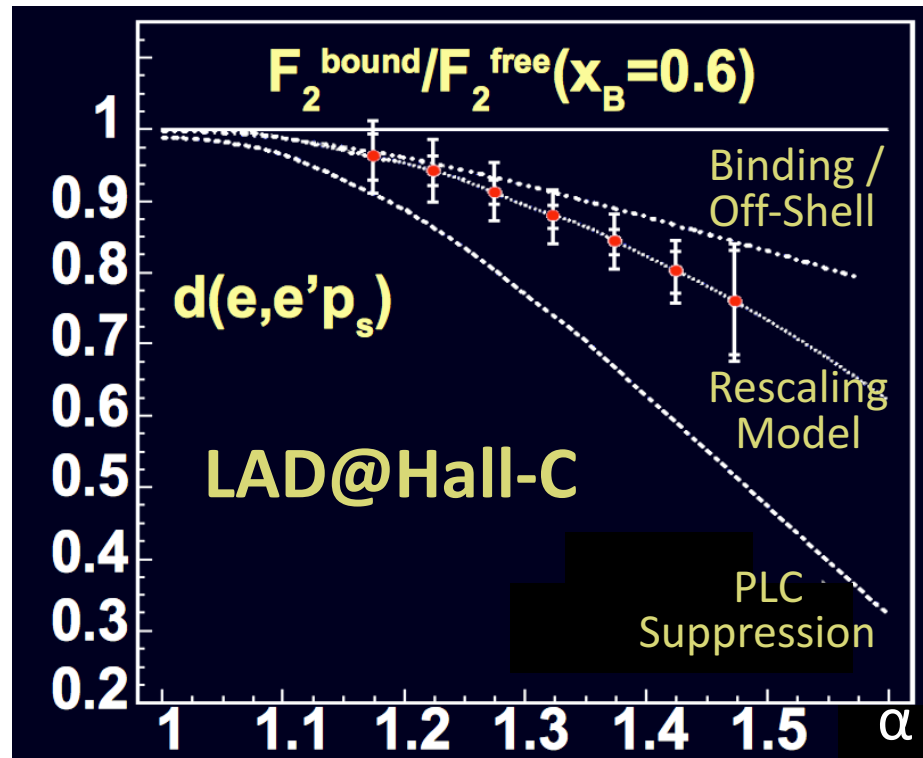
Internal structure of SRC nucleons?

Focus on the deuteron:

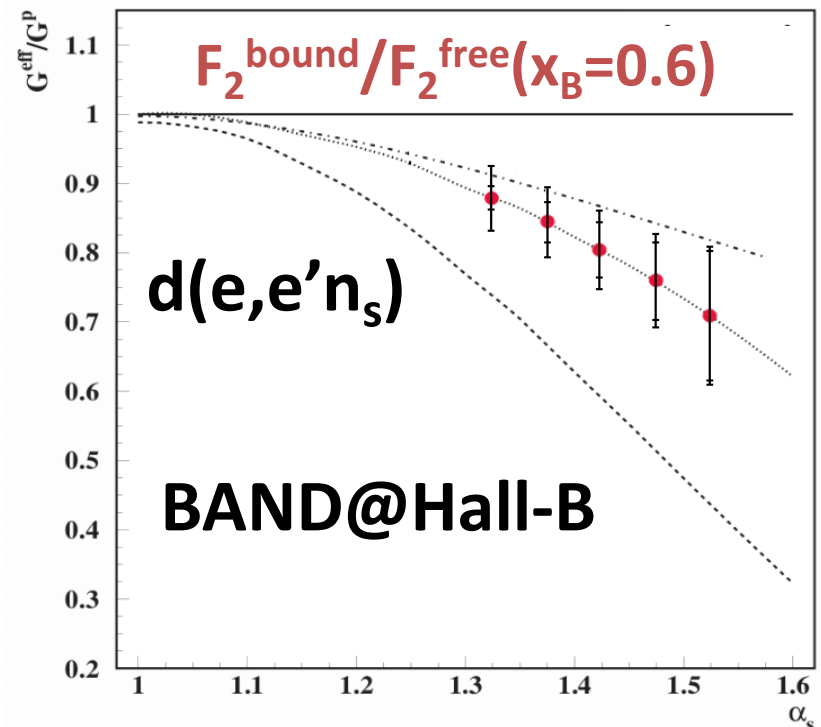
(2) Infer its momentum from the recoil partner.



(1) Perform DIS off forward going nucleon.

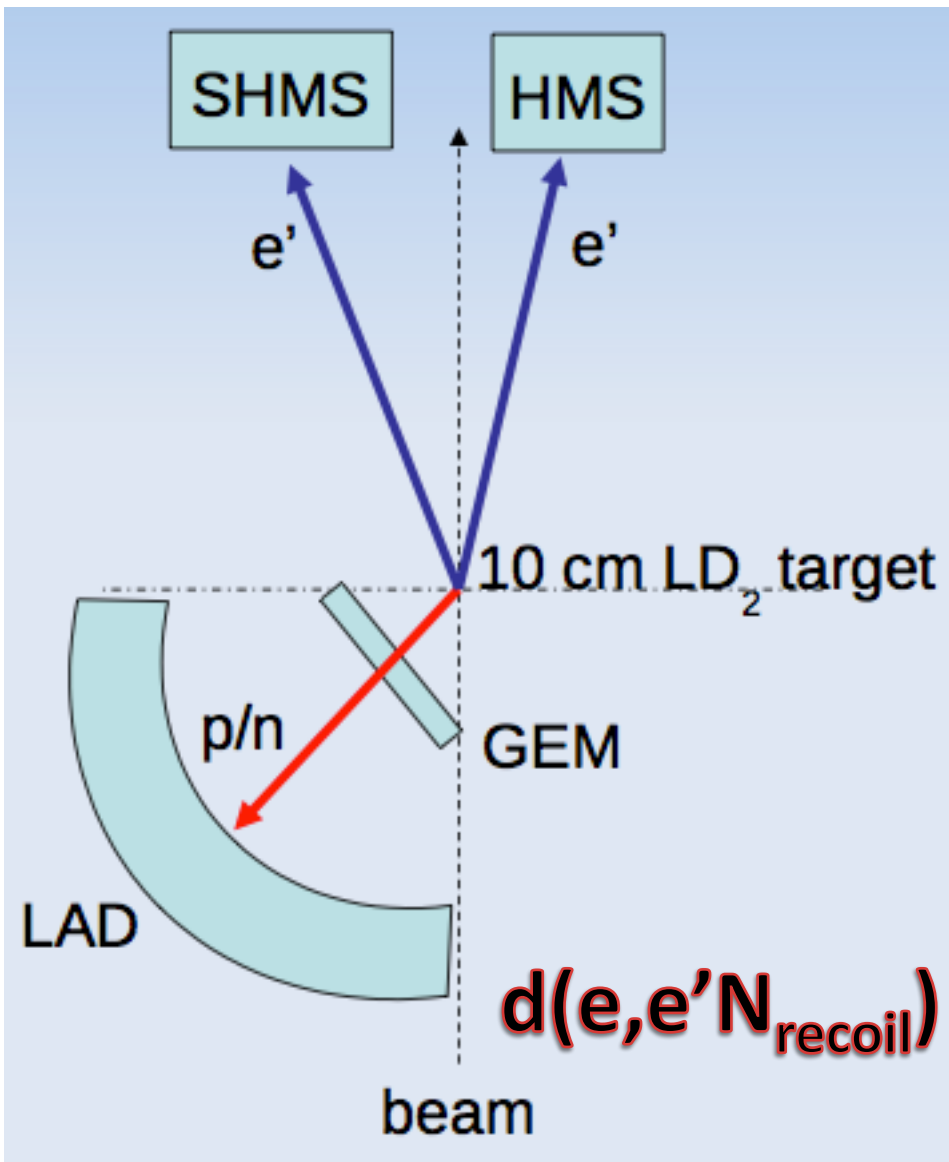


Melnitchouk et al., Z. Phys. A **359**, 99-109 (1997)





Fixed Target Concept...



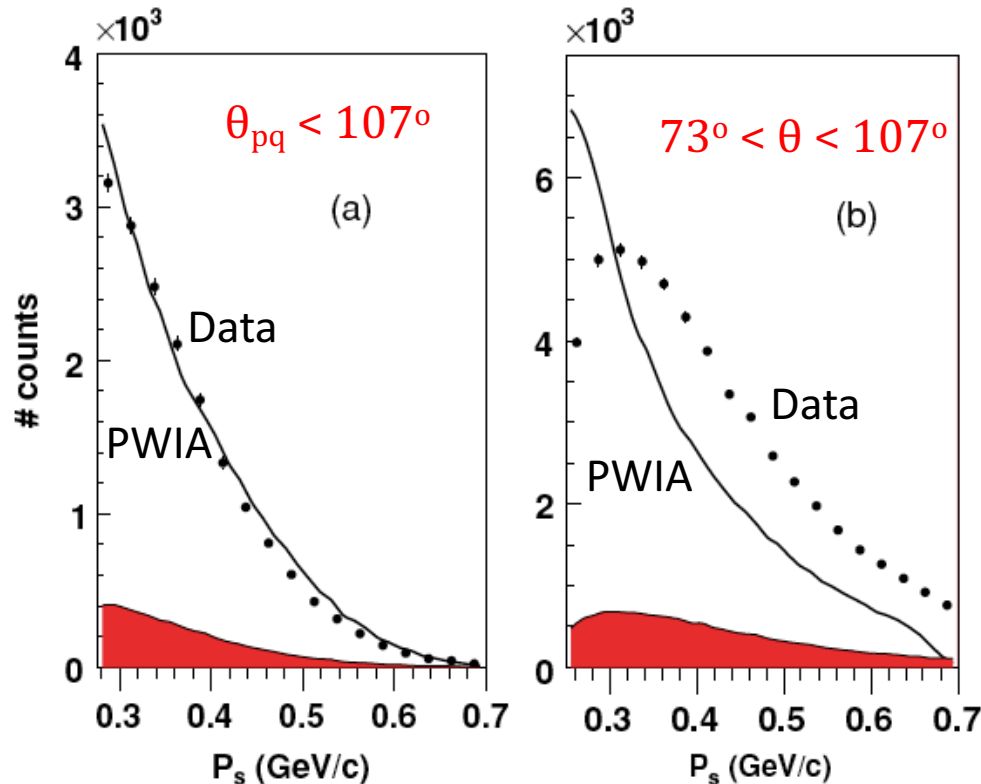
- High resolution spectrometers for $d(e, e')$ measurement in DIS kinematics
- Large acceptance recoil proton \ neutron detector
- Long target + GEM detector – reduce random coincidence



Backward Kinematics:

Minimize Re-Scattering

$D(e, e' p_s)$



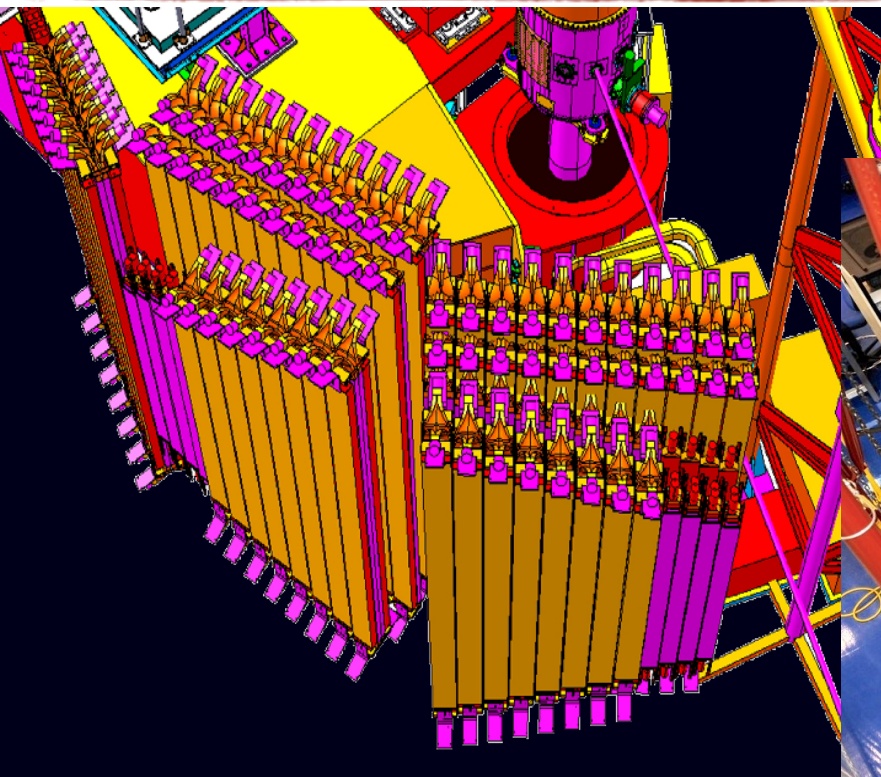
A. V. Klimenko *et al.*, PRC 73, 035212 (2006)

FSI:

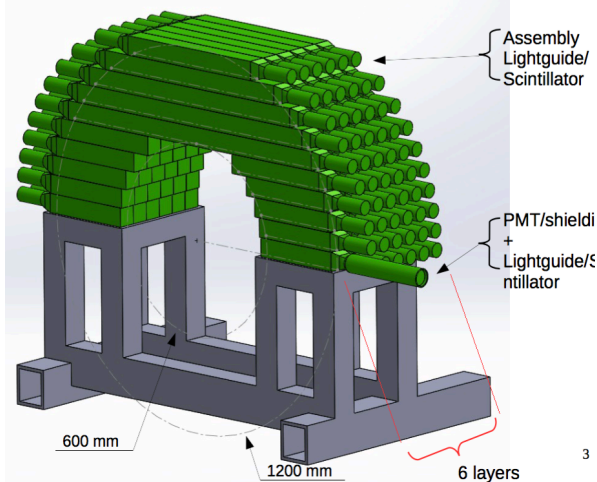
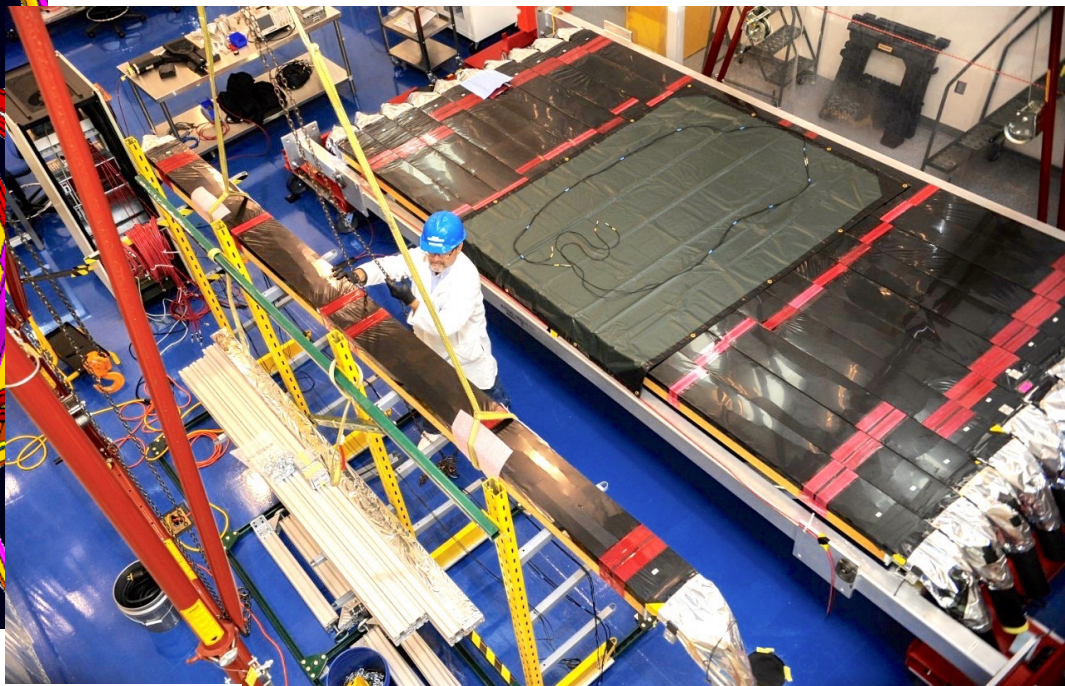
- Decrease with Q^2
- Increase with W'
- Not sensitive to x'
- Small for $\theta_{pq} > 107^\circ$



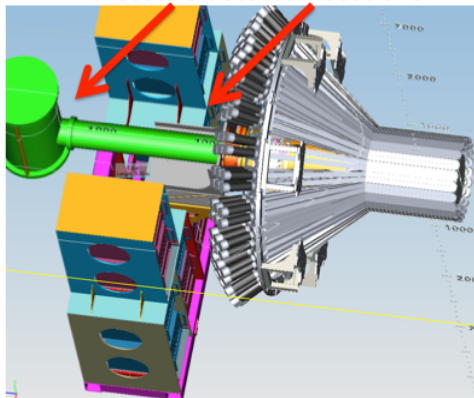
Next Generation Experiments



Large Acceptance
Detector (LAD@Hall-C)



Possible detector locations



Backward Angle Neutron
Detector (BAND@Hall-B)

R&D @ MIT /
Construction @ BATES

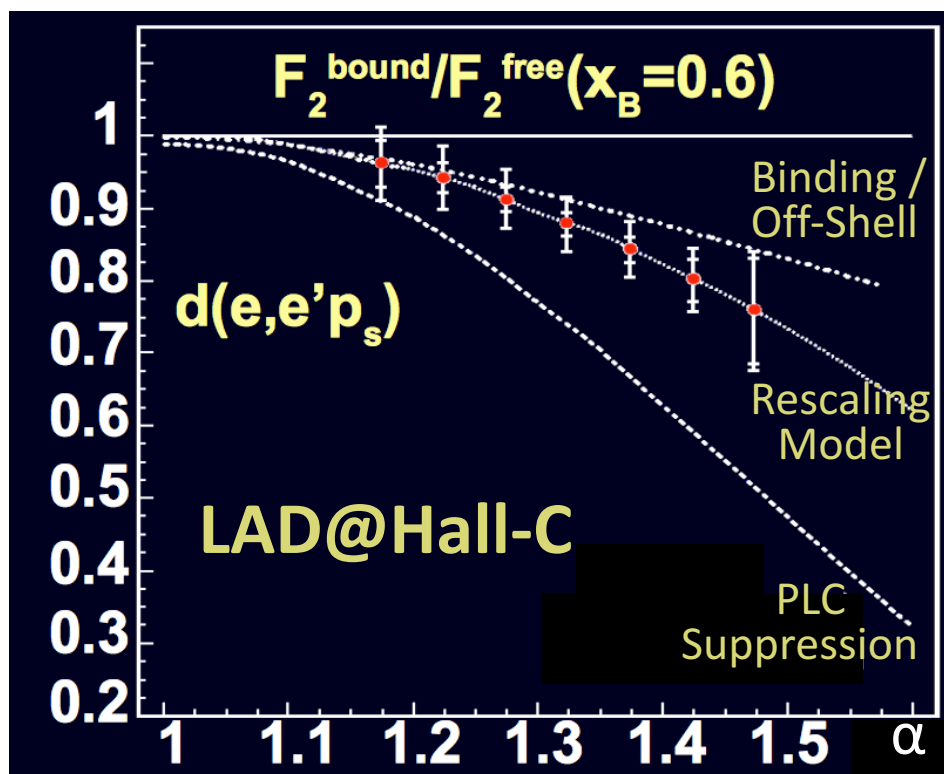


Kinematics and Uncertainties

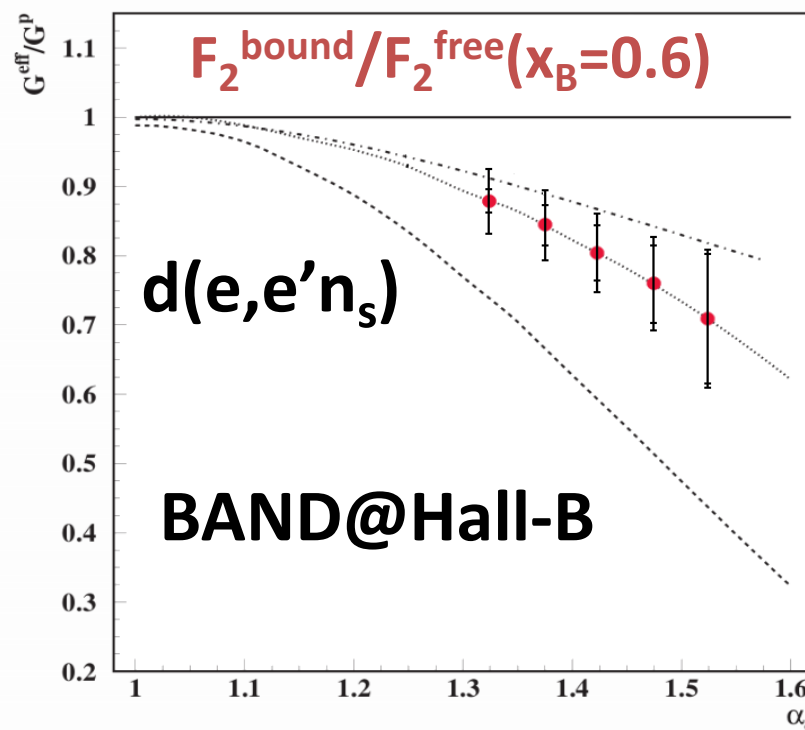


- Tagging allows to extract the structure function in the nucleon reference frame: $x' = \frac{Q^2}{2(\bar{q} \cdot \bar{p})}$
- Expected coverage: $x' \sim 0.3$ & $0.45(0.5) < x' < 0.55(0.7)$ @

$$W^2 > 4 \text{ [GeV/c]}^2$$



Melnitchouk et al., Z. Phys. A **359**, 99-109 (1997)





The Correlations group



- MIT (Or Hen):



Barak Schmookler



Reynier Torres



Efrain Segarra



Afroditi Papadopoulou



Axel Schmidt



George Laskaris



Maria Patsyuk



Taofeng Wang (*visiting scientist)

- TAU (Eli Piassetzky):



Erez Cohen



Meytal Duer



Igor Korover



Adi Ashkenazy

- ODU (Larry Weinstein):



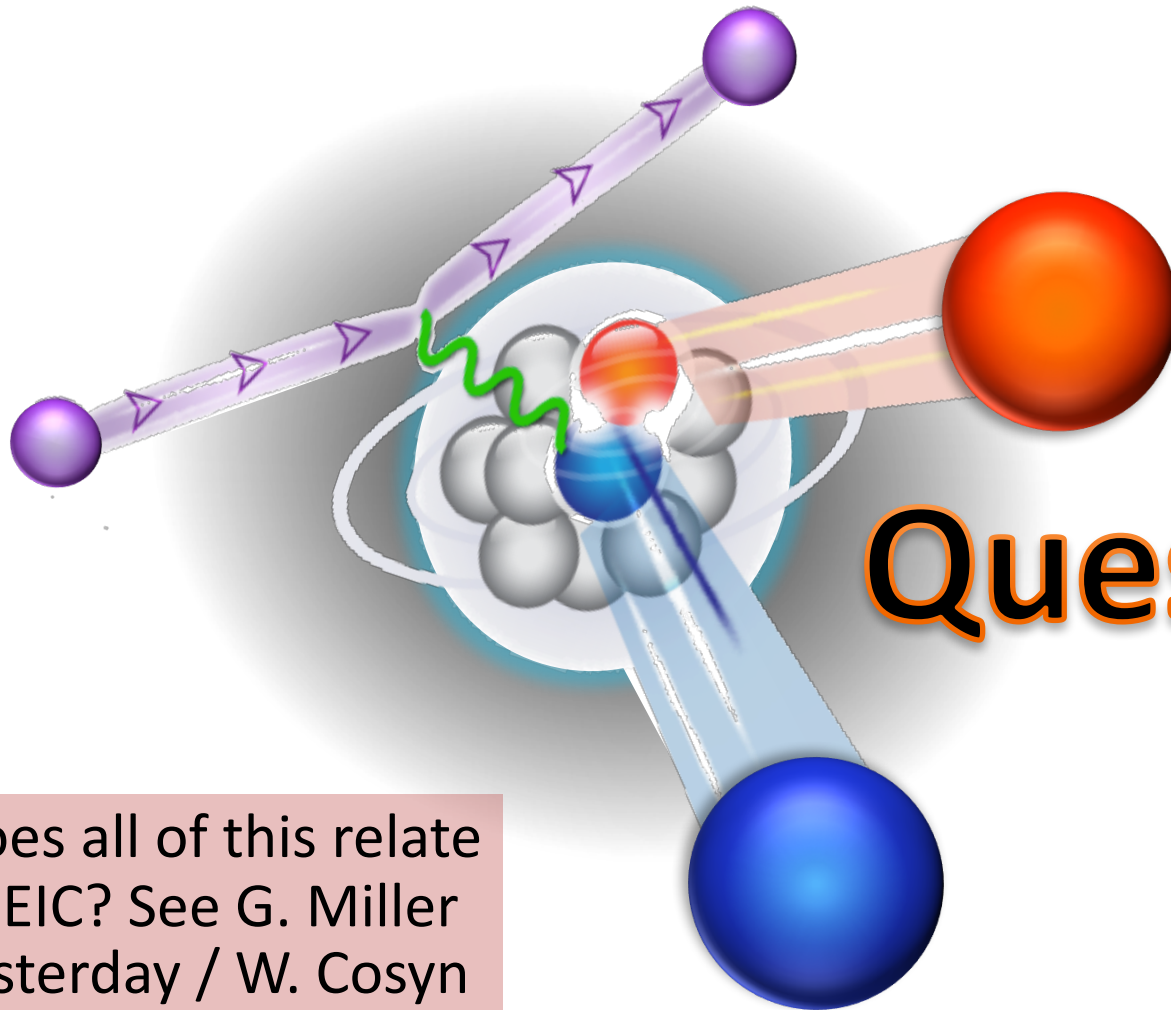
Mariana Khachatryan



Florian Hauenstein

- Theory Collaborators (lots!)

Thank You!



Questions?

How does all of this relate to the EIC? See G. Miller talk yesterday / W. Cosyn talk tomorrow.